

The 4th Trilateral International Workshop on Energetic Particle Physics

Simulation of β -induced Alfvén eigenmode instabilities and mode transition for HL-3 hybrid scenario

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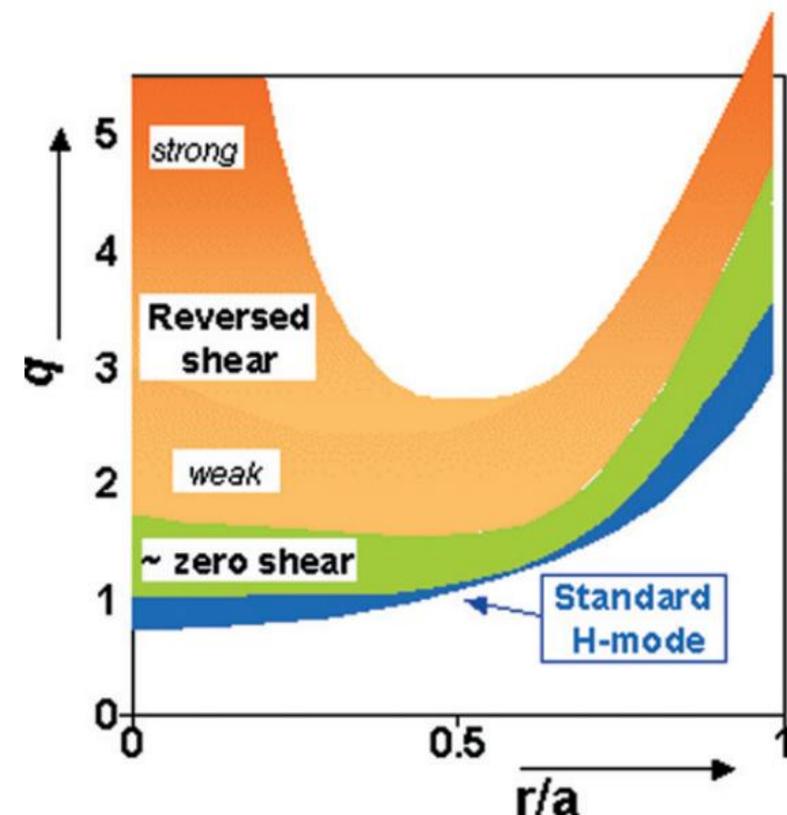


- **Motivation**
- **HL-3 hybrid scenario and simulation method**
- **Simulation results**
 - Multiple-n (n=1-5) simulation**
 - Parametric ($\beta_{EP} / E_{inj} / q$) dependence on the n=2 mode**
 - Nonlinear simulation of n=1 and n=2**
- **Summary**



■ Hybrid scenario / Hybrid scenario

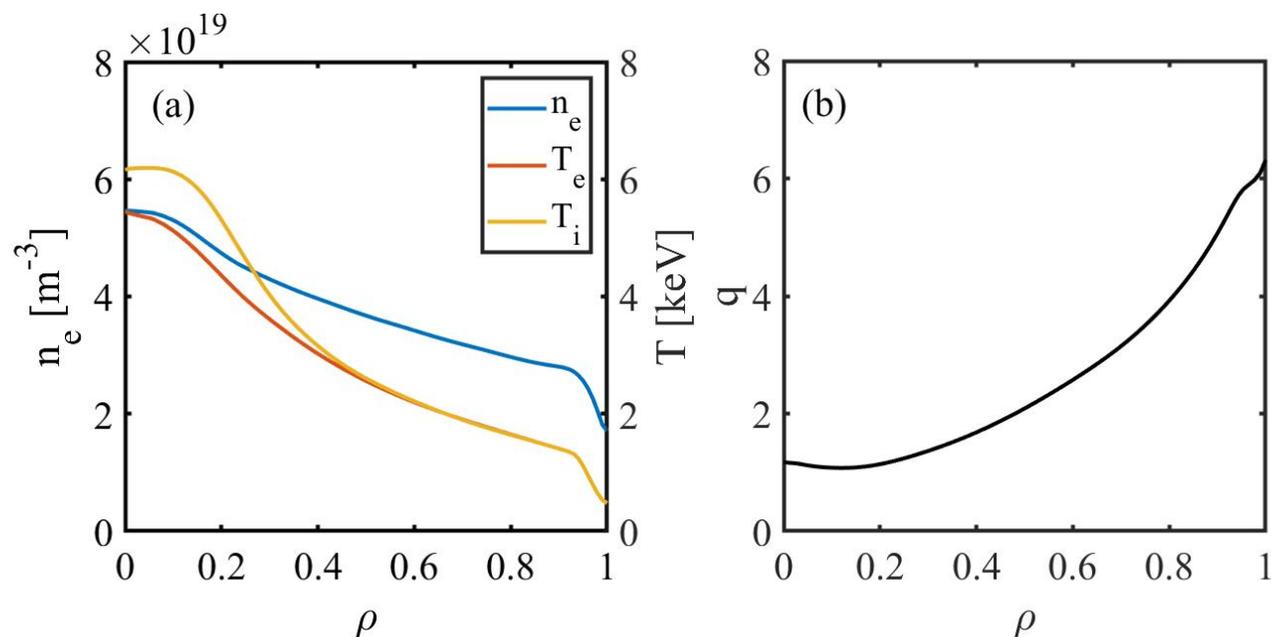
- Ohmic fraction $\neq 0$
- Core instabilities due to the shape of q -profile
- Higher fusion power and fusion gain
- Alfvén eigenmodes
- HL-3 tokamak 1MA H-mode



The range of q -profiles for the ITER

HL-3 hybrid scenario

- $B=2.2\text{T}$, $I_p=1\text{MA}$, $R=1.78\text{m}$, $a=0.6\text{m}$
- Heating: Ohmic+NBI+ECRH
- Off-axis ECCD results in weak reverse shear on-axis
- $q_{min}>1$ to avoid sawtooth
- Equilibrium and plasma profiles are generated by the kEFIT workflow through OMFIT integrated simulation framework
- The instabilities are always excited in the flat shear region



The bulk plasma is described by nonlinear MHD equations:

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho v) + v_n \nabla^2 (\rho - \rho_{eq}) \quad (1)$$

$$\rho \frac{\partial v}{\partial t} = -\rho \omega \times v - \rho \nabla \left(\frac{v^2}{2} \right) - \nabla p + (j - j_h) \times B - \nabla \times (v \rho \omega) + \frac{4}{3} (v \rho \nabla \cdot v) \quad (2)$$

$$\frac{\partial B}{\partial t} = -\nabla \times E \quad (3)$$

$$\frac{\partial p}{\partial t} = -\nabla \cdot (p v) - (\Gamma - 1) p \nabla \cdot v + (\Gamma - 1) \left[v \rho \omega^2 + \frac{4}{3} v \rho (\nabla \cdot v)^2 + \eta j \cdot (j - j_{eq}) \right] + \chi \Delta (p - p_{eq}) \quad (4)$$

$$E = -v \times B + \eta (j - j_{eq}) \quad (5)$$

$$\omega = \nabla \times v \quad (6)$$

$$j = \frac{1}{\mu_0} \nabla \times B \quad (7)$$

EP current density:

$$j_h = \int (v_{\parallel}^* + v_B) Z_h e f d^3 v - \nabla \times \int \mu b f d^3 v \quad (8)$$

EP slowing down distribution:

$$f_0(\psi, v, \Lambda) = C \exp\left(-\frac{\psi}{\Delta \psi}\right) \frac{1}{v^3 + v_c^3} \times \frac{1}{2} \operatorname{erfc}\left(\frac{v - v_0}{\Delta v}\right) \times \exp\left(-\frac{(\Lambda - \Lambda_0)^2}{\Delta \Lambda^2}\right) \quad (9)$$

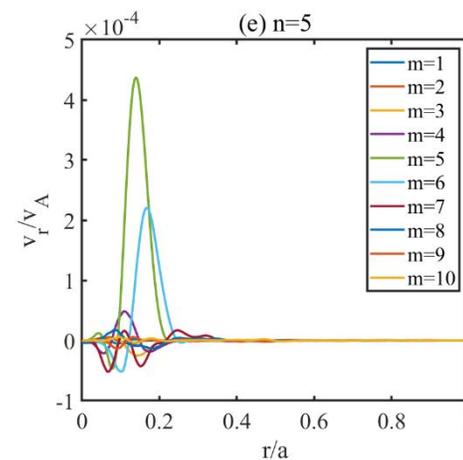
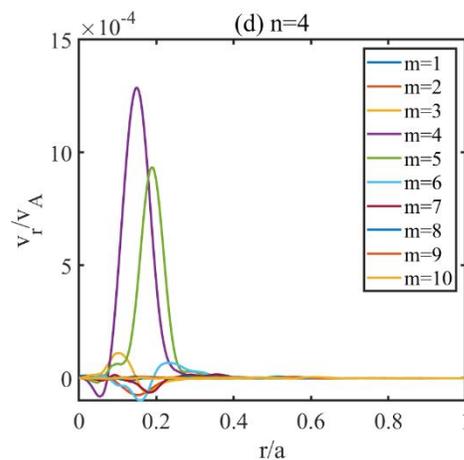
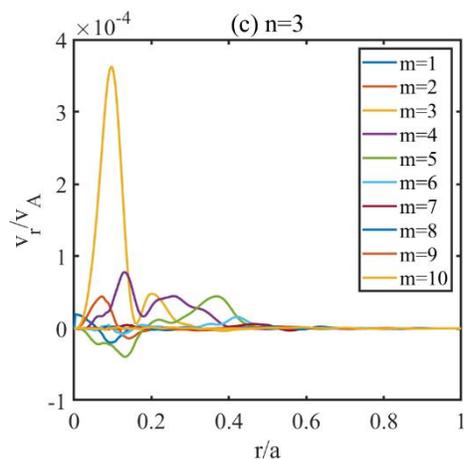
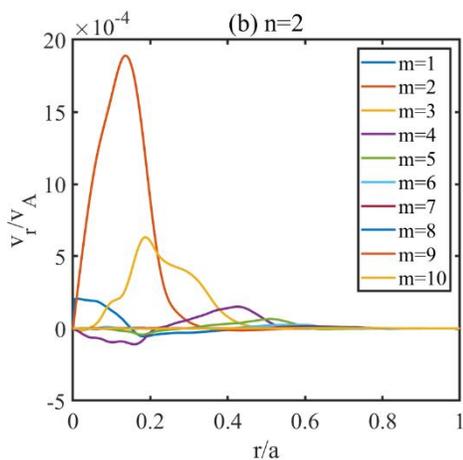
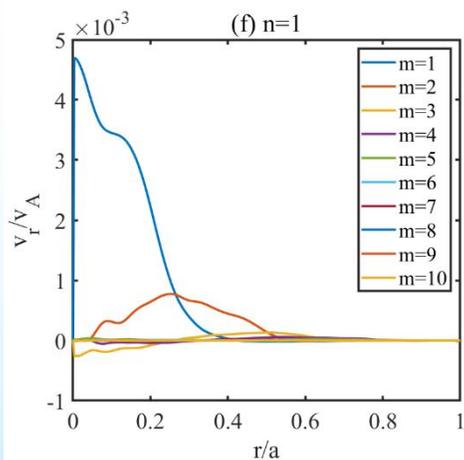
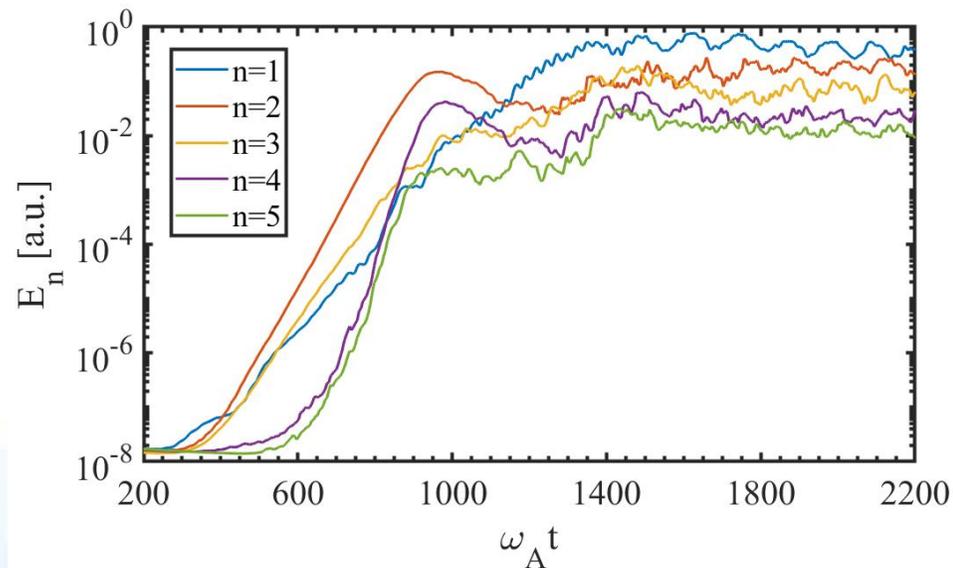


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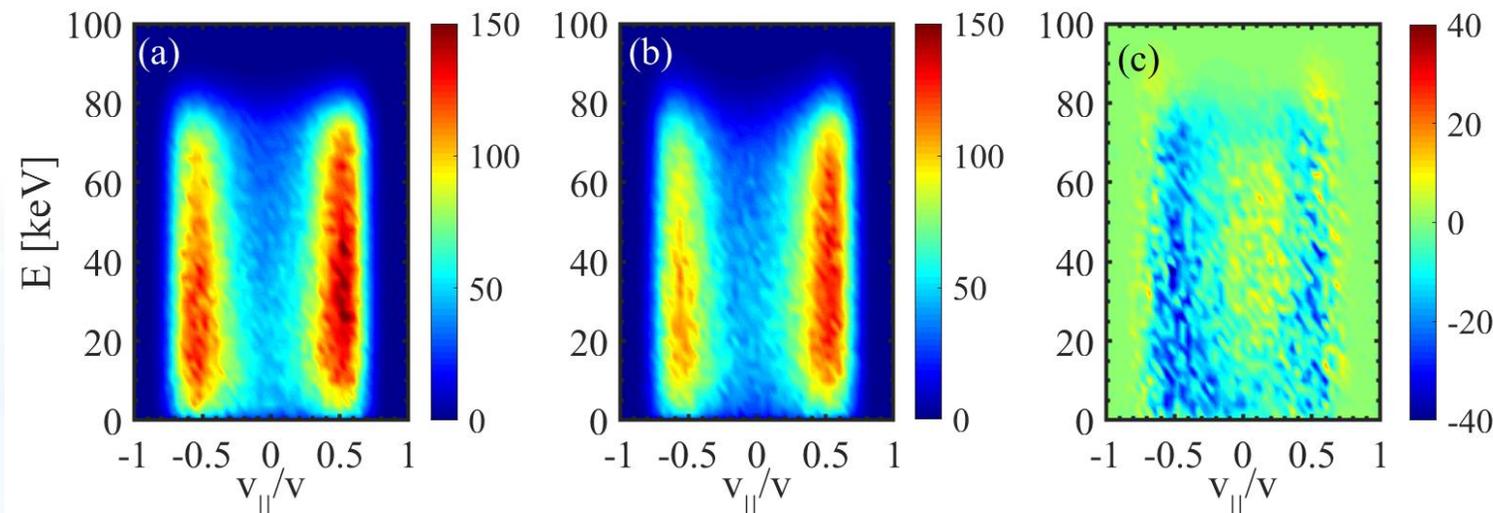
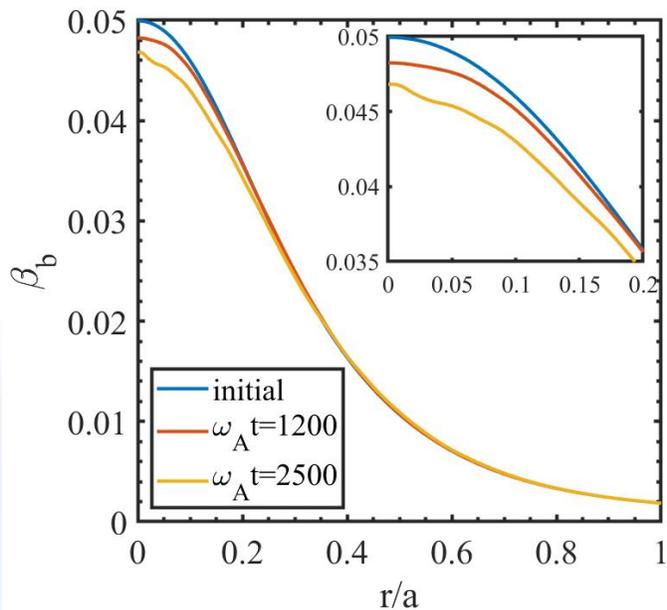


Multiple-n simulation to determine the most unstable mode

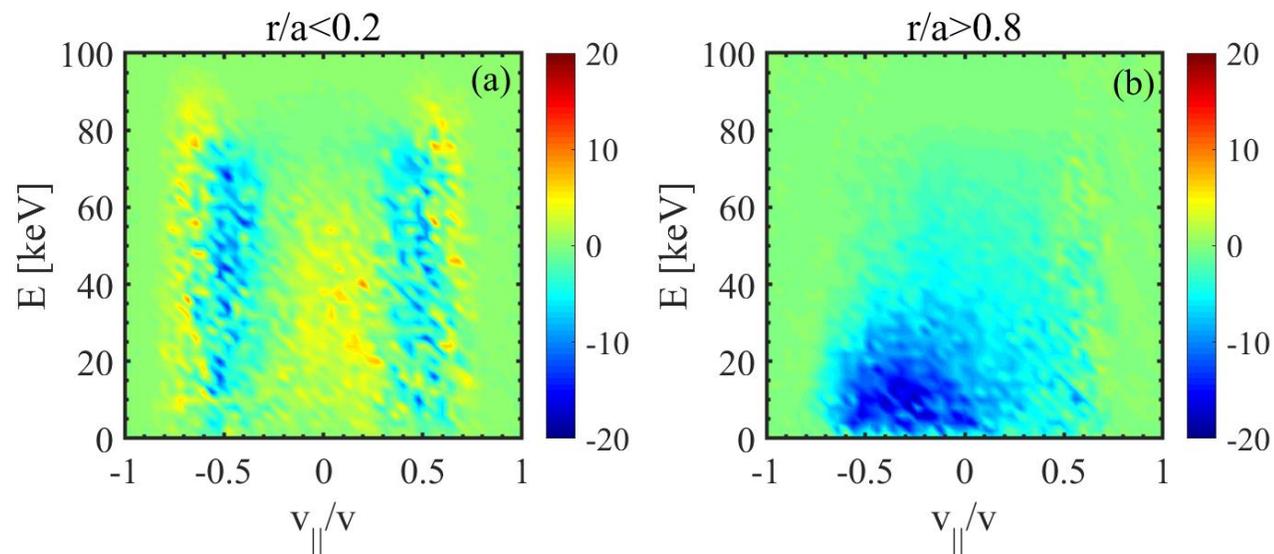
- Toroidal mode number: $n=1-5$
- Assumption: $\beta_{EP} \sim \beta_{th}$, $E_{inj}=80\text{keV}$
- The $n=2$ mode (BAE) is firstly driven unstable
- The saturated amplitude of the $n=1$ mode (fishbone) is the largest
- The saturated amplitude of $n=3-5$ is much lower



Energetic particle redistribution



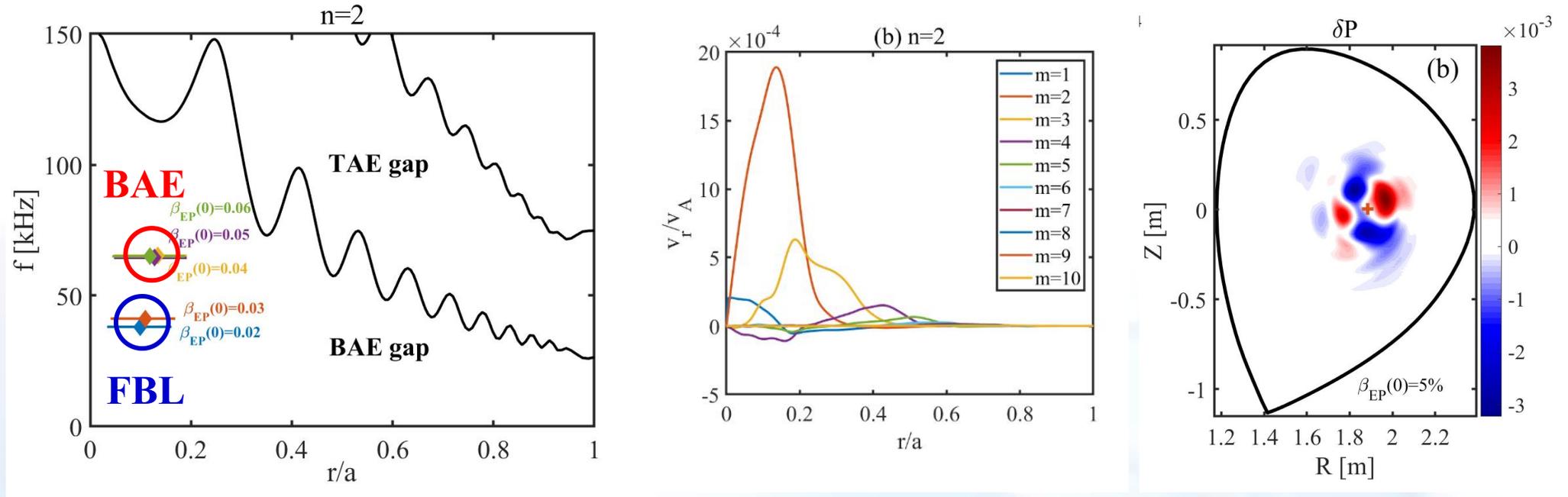
- EP profile decrease is insignificant near the axis with the reduction of $\beta_{EP}(0) \sim 3.42\%$.
- The reduction of counter-moving particles is more significant than co-moving particles
- Core($r/a < 0.2$): redistribution
- Edge($r/a > 0.8$): loss



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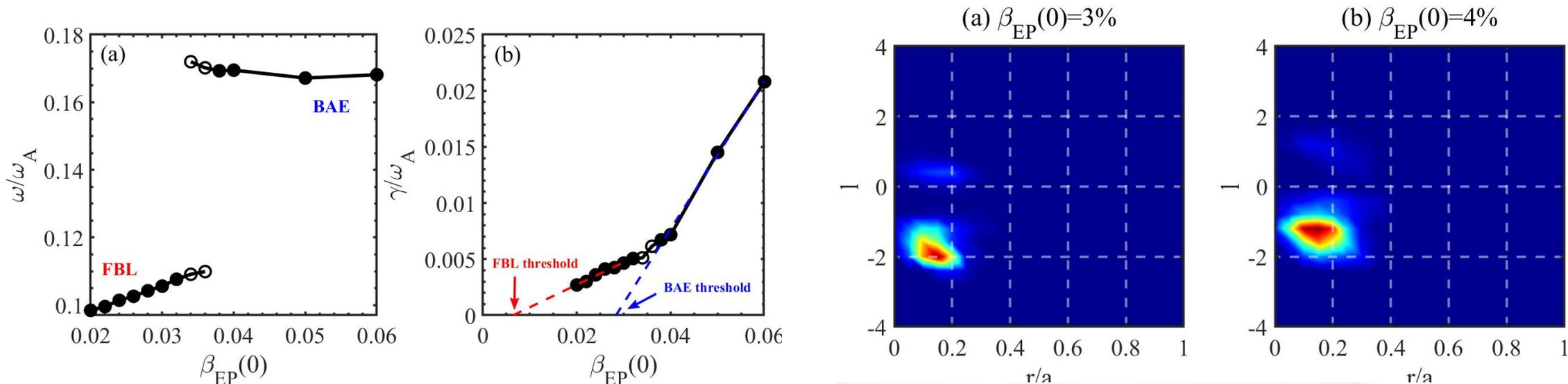
The effect of EP pressure on the $n=2$ mode



- Raising EP pressure, the most unstable mode becomes a beta-induced Alfvén eigenmode (BAE) from a fishbone-like mode (FBL)
- The $m=2$ harmonic is dominant for the two modes
- The mode structure of the FBL and BAE are similar, locating at $q=1$ flux surface

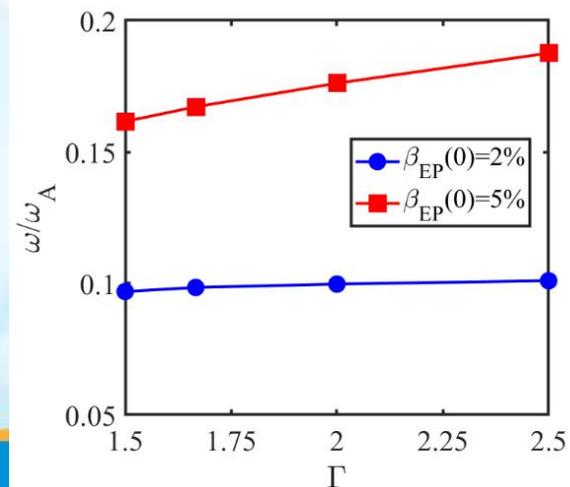


Differences between the two n=2 modes

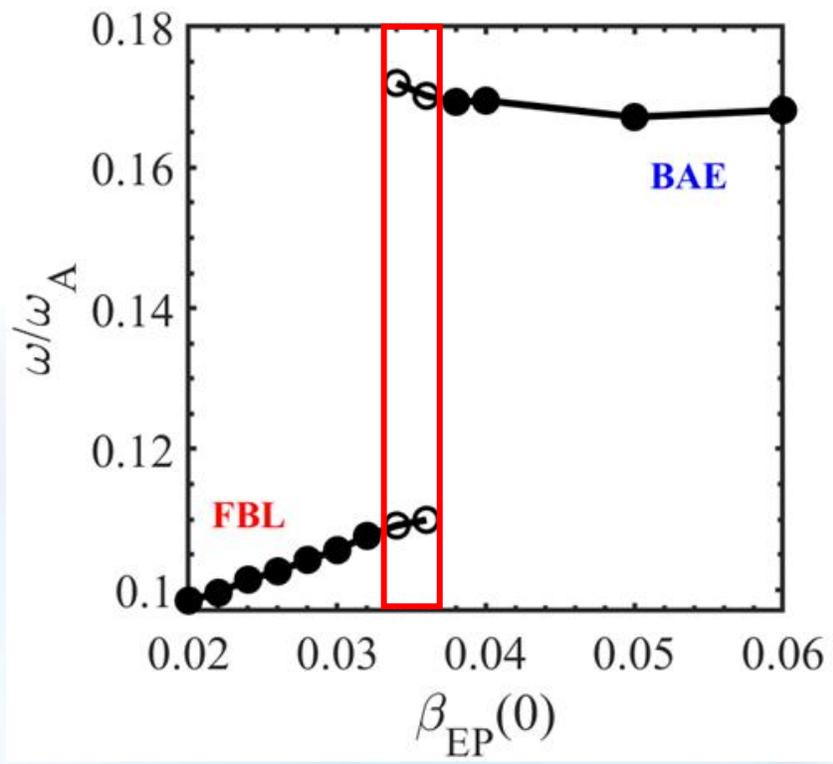


wave-particle resonant condition: $\omega - l\omega_\theta - n\omega_\phi = 0$

	FBL	BAE
Frequency (f)	45kHz	60kHz
Linear growth rate (γ)	Small	Large
Adiabatic constant (Γ)	No effect	Linear
Resonant number (l)	-2	-1



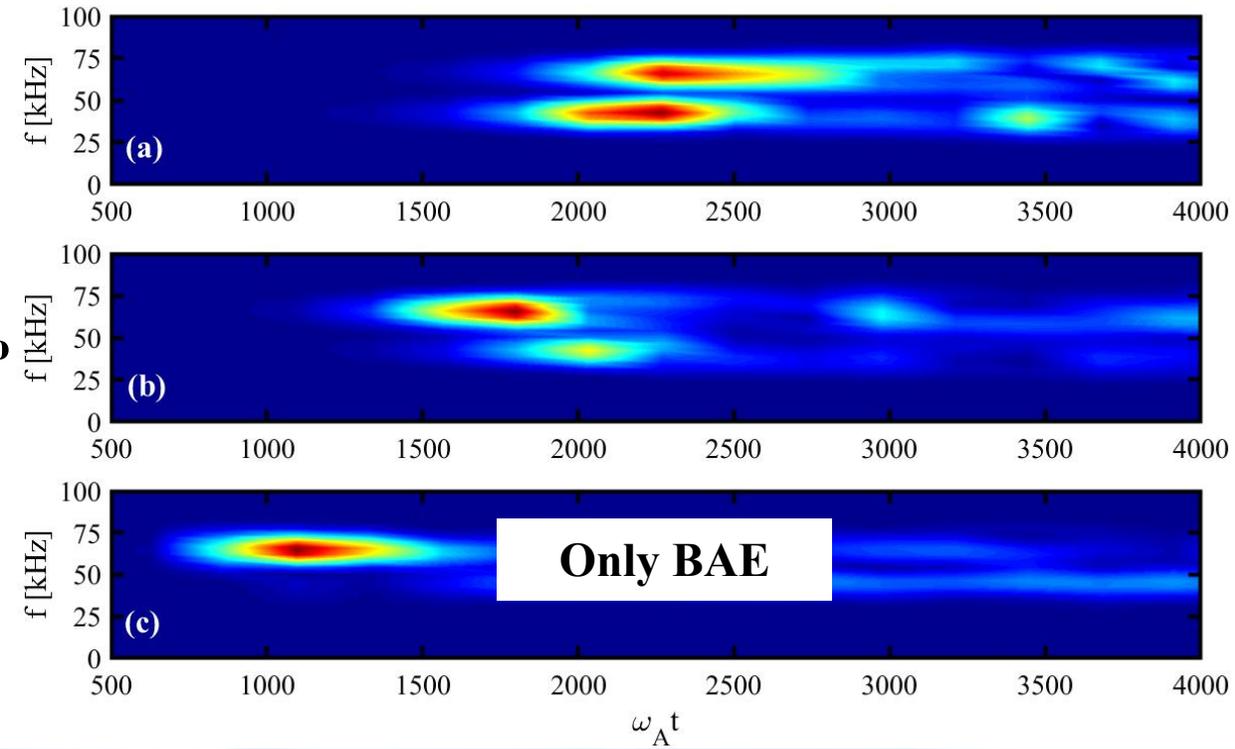
Co-existence of the two n=2 modes



$\beta_{EP}=3.4\%$

$\beta_{EP}=3.6\%$

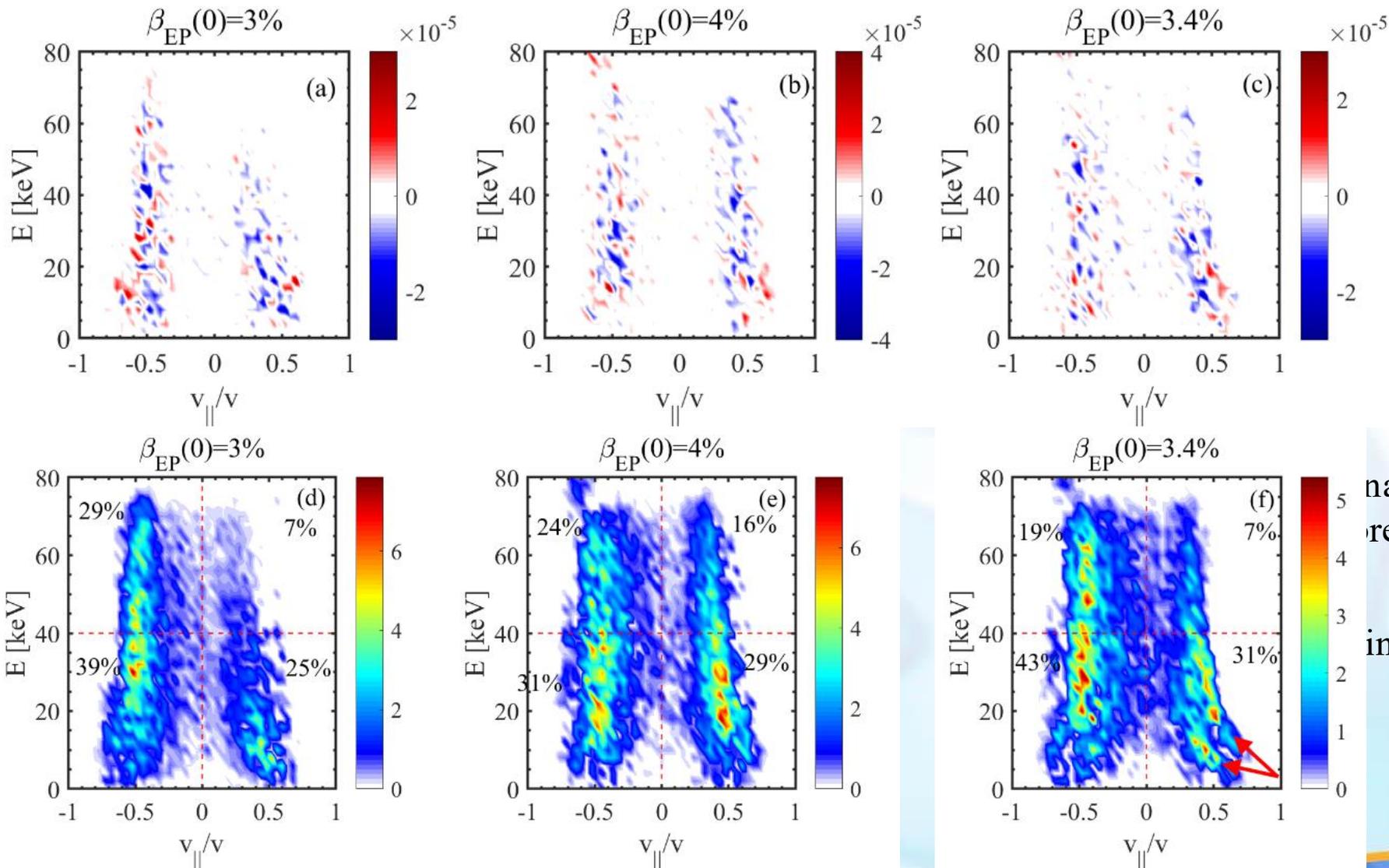
$\beta_{EP}=5\%$



The two n=2 modes (FBL and BAE) can co-exist within $\beta_{EP}=3.4\%-3.6\%$.
 The FBL or BAE excited firstly is related to the EP pressure



Perturbed distribution δf and distribution f



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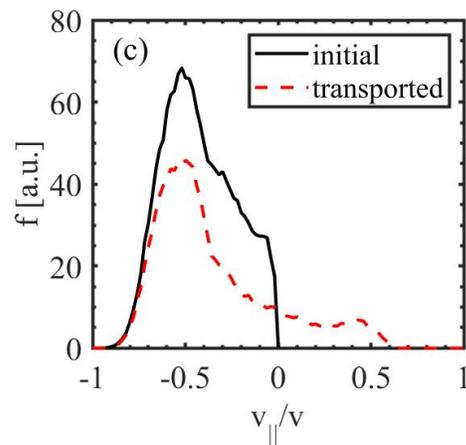
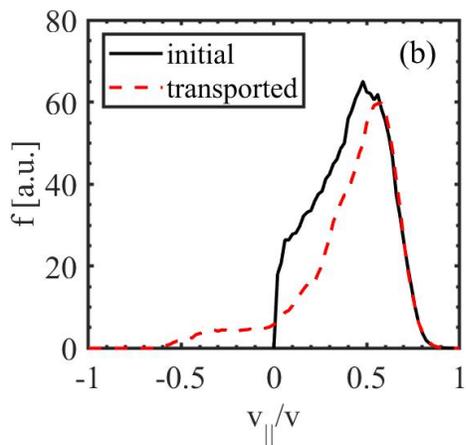
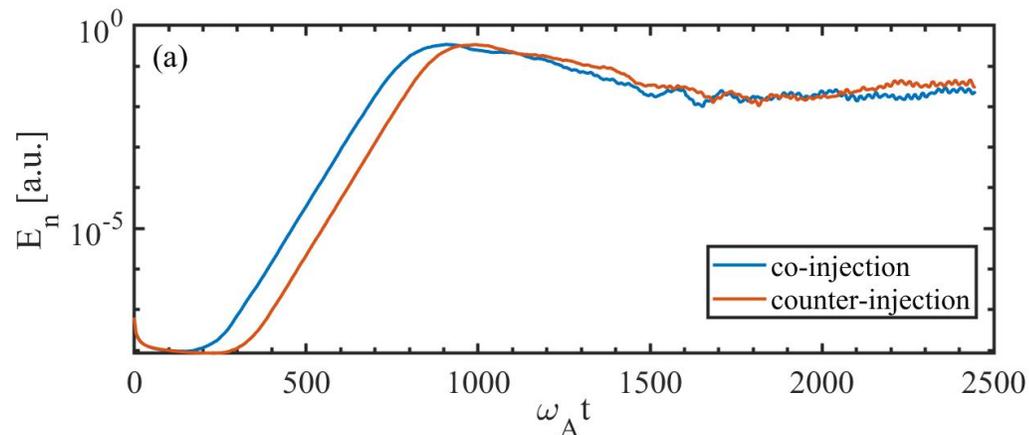
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Two mode co-exist



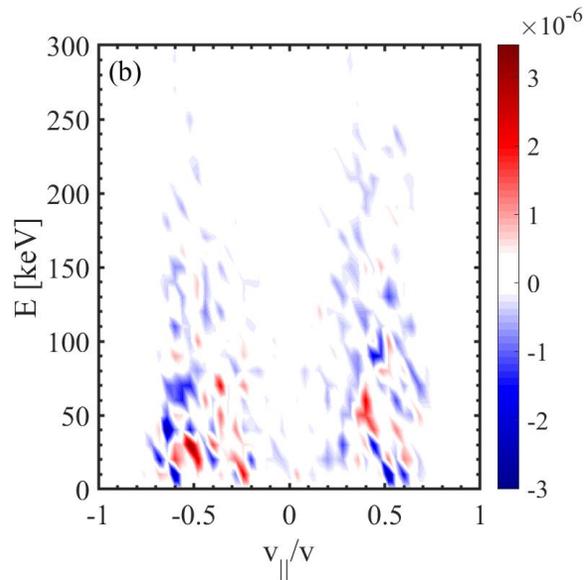
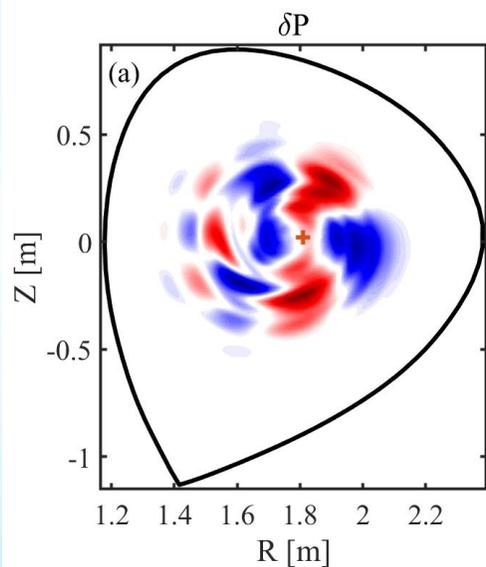
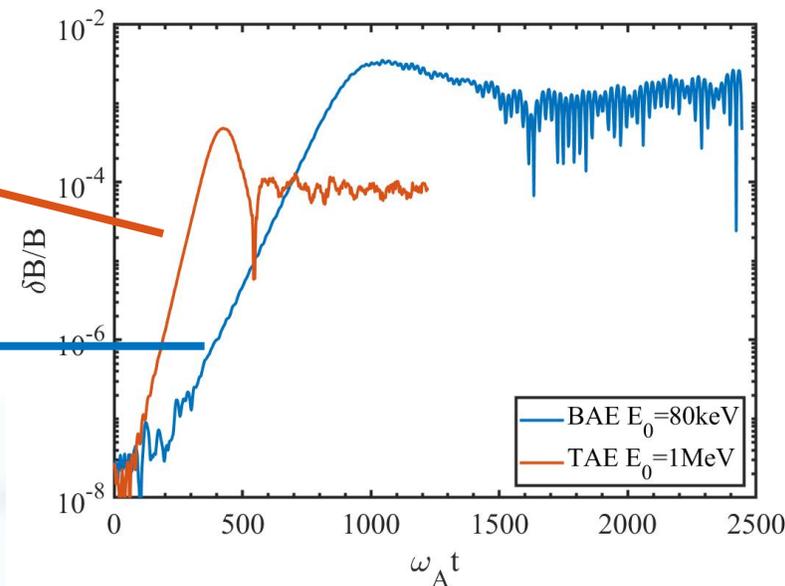
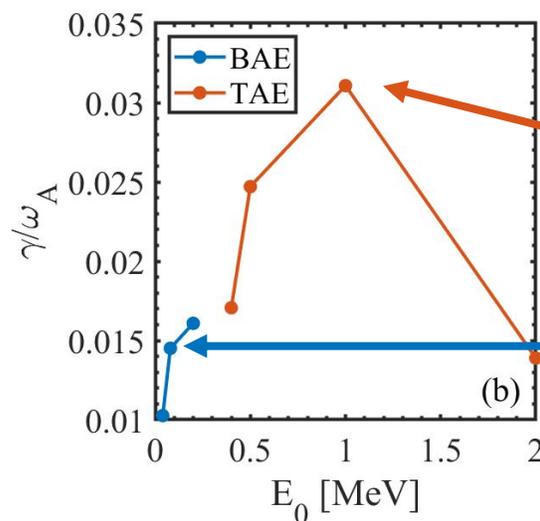
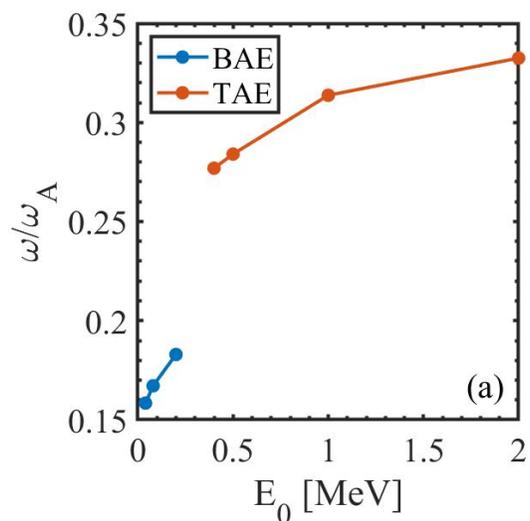
The effect of injected direction on the EP distribution



- Select two extreme ideal situations: the deposited distribution exists only in positive (co-injection) and negative (counter-injection) v_{\parallel}/v region.
- The peak for the counter-injection case showing a much more significant reduction than co-injection. (e.g. 50keV)



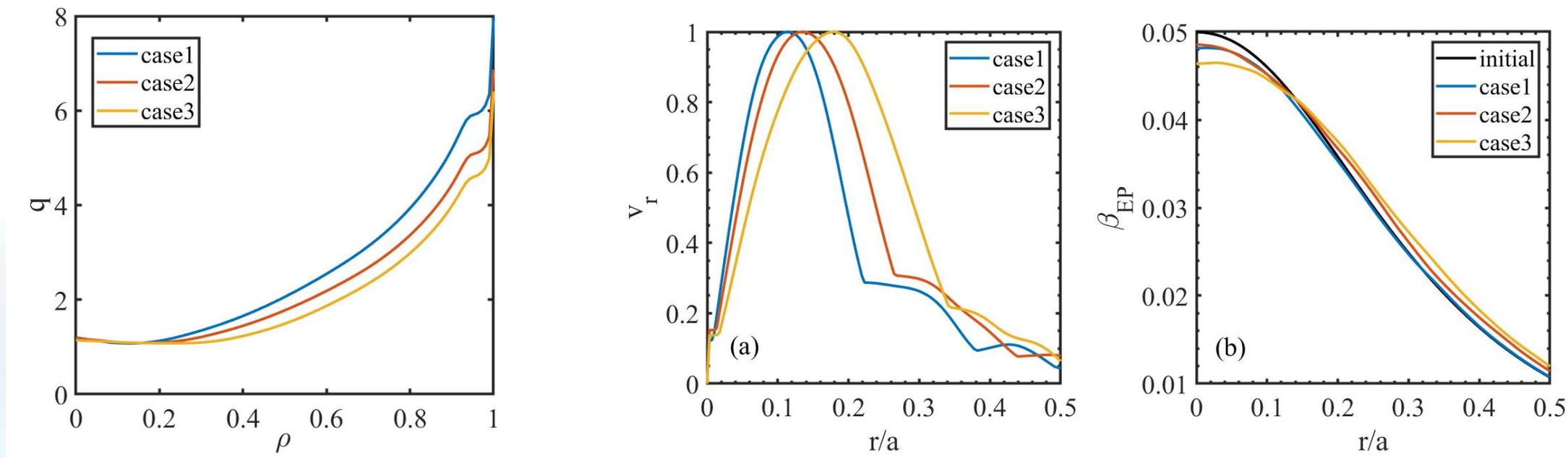
The effect of injected energy on the n=2 mode



- Raising EP pressure, the most unstable mode becomes a toroidal Alfvén eigenmode (TAE) from a beta-induced Alfvén eigenmode (BAE)
- Compared different injected energy, although the linear growth rate of TAE is larger, its saturated amplitude is lower



The effect of q -profile on the BAE



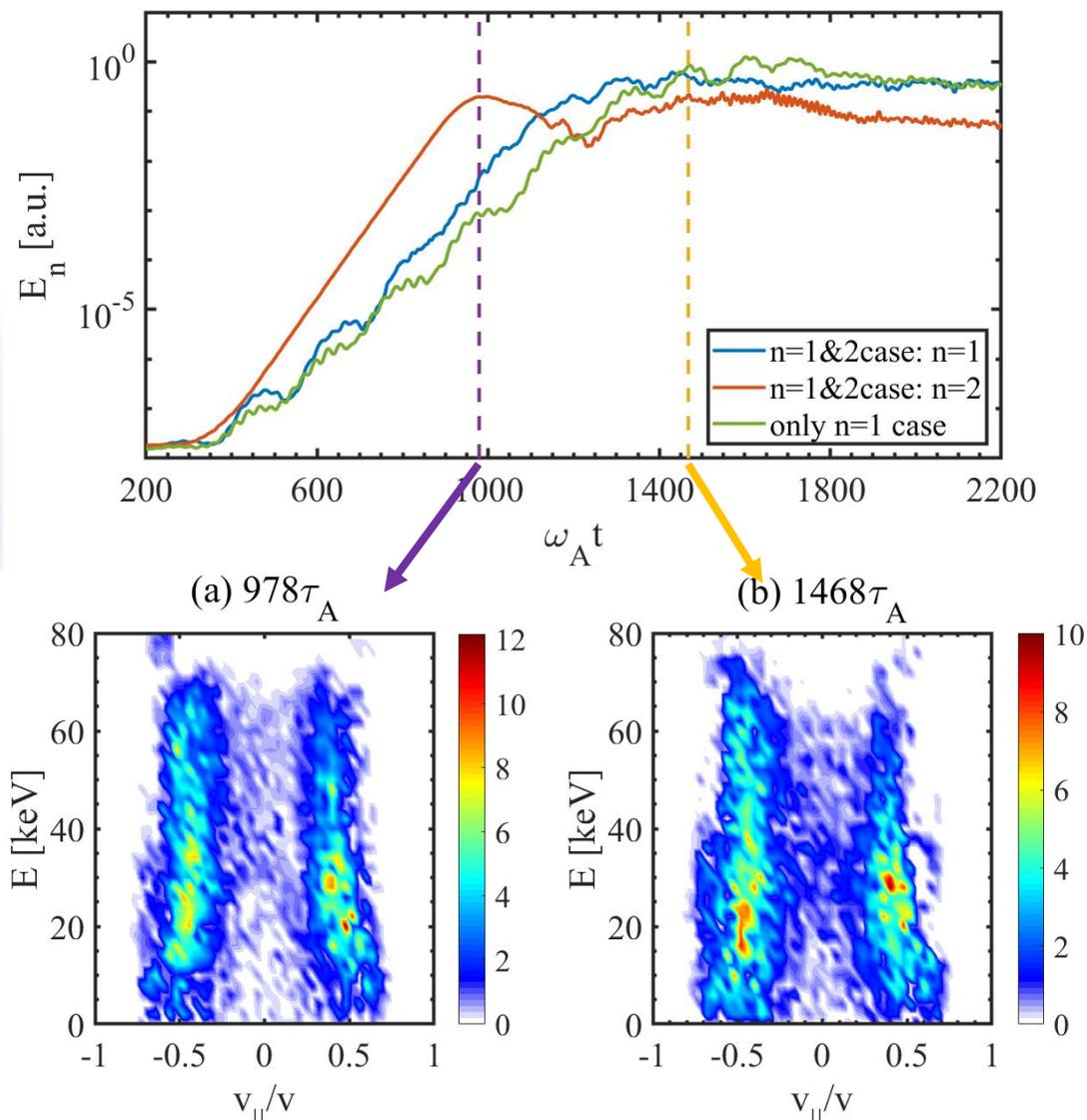
- The EFIT code is used to extend the flat shear region with fixed q_{\min} (Thermal pressure keeps fixed)
- The mode structure is the envelope of $m=1-5$ harmonics, and normalized to the peak value
- With extension of the flat shear region, the width of mode structure increases (The full width at half maxima increases from $0.16a$ to $0.23a$), resulting in more EP transport.



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Simulation with only considering n=1 & 2 modes



- Case1: only n=1
- Case2: n=1 & 2

- For the two cases, the saturated energy of n=1 mode are similar, but the linear growth rate of n=1 mode is slightly larger when considering the n=2 mode
- EP acts as an intermediate to transfer energy from n=2 mode to n=1 mode.



Summary

- HL-3 heating & current drive:

Ohmic

ECRH: 2MW

NBI: 2MW→4MW

- For HL-3 hybrid scenario (1MA): low n mode

- For the n=2 mode

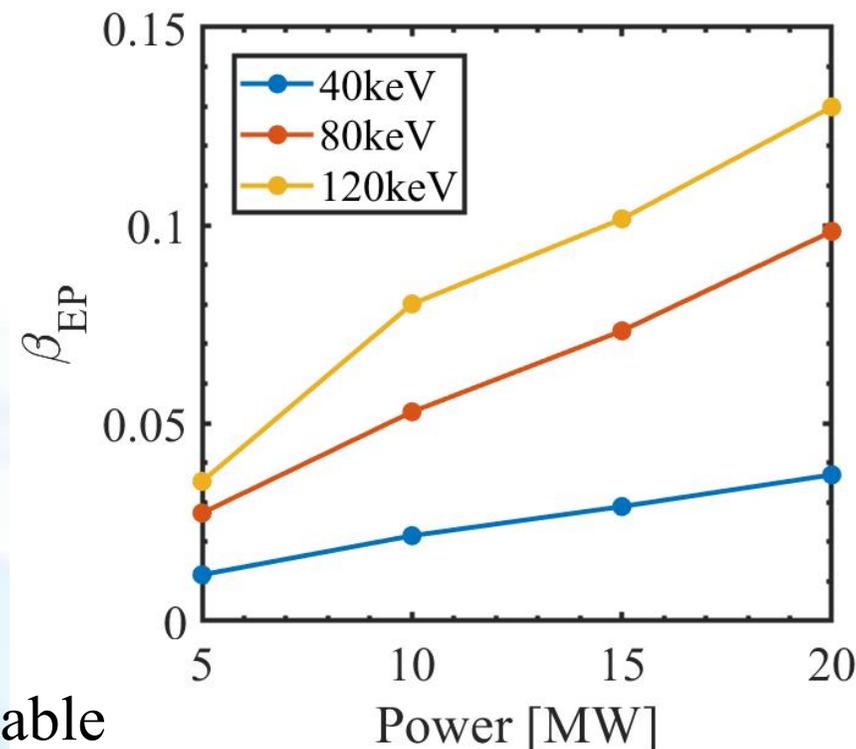
Raising EP pressure: FBL→BAE

Raising injected energy : BAE → TAE

Broadening flat shear region, BAE becomes more unstable

- EP acts as an intermediate to transfer energy between modes

- More EPs will be observed in HL-3 device by increasing heating power



Thanks for your attention

