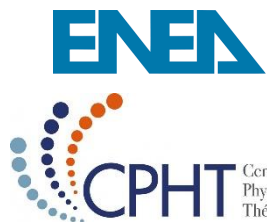




Multiscale Synergies in Fusion Plasmas: Turbulence, MHD Modes, Fast Particles, and Alfvénic Instabilities

A. Mishchenko, F. Antlitz, A. Biancalani, A. Bottino, M. Borchardt, M. Falessi, T. Hayward-Schneider, R. Kleiber, A. Könies, E. Lanti, P. Lauber, Z. Lu, Y. Narbutt, C. Nührenberg, E. Poli, J. Riemann, B. Rofman, J. N. Sama, C. Slaby, L. Villard, X. Wang, F. Widmer, F. Zonca



This work has been carried out within the framework of the EUROfusion Consortium, funded by the European Union via the Euratom Research and Training Programme (Grant Agreement No 101052200 — EUROfusion). Views and opinions expressed are however those of the author(s) only and do not necessarily reflect those of the European Union or the European Commission. Neither the European Union nor the European Commission can be held responsible for them.

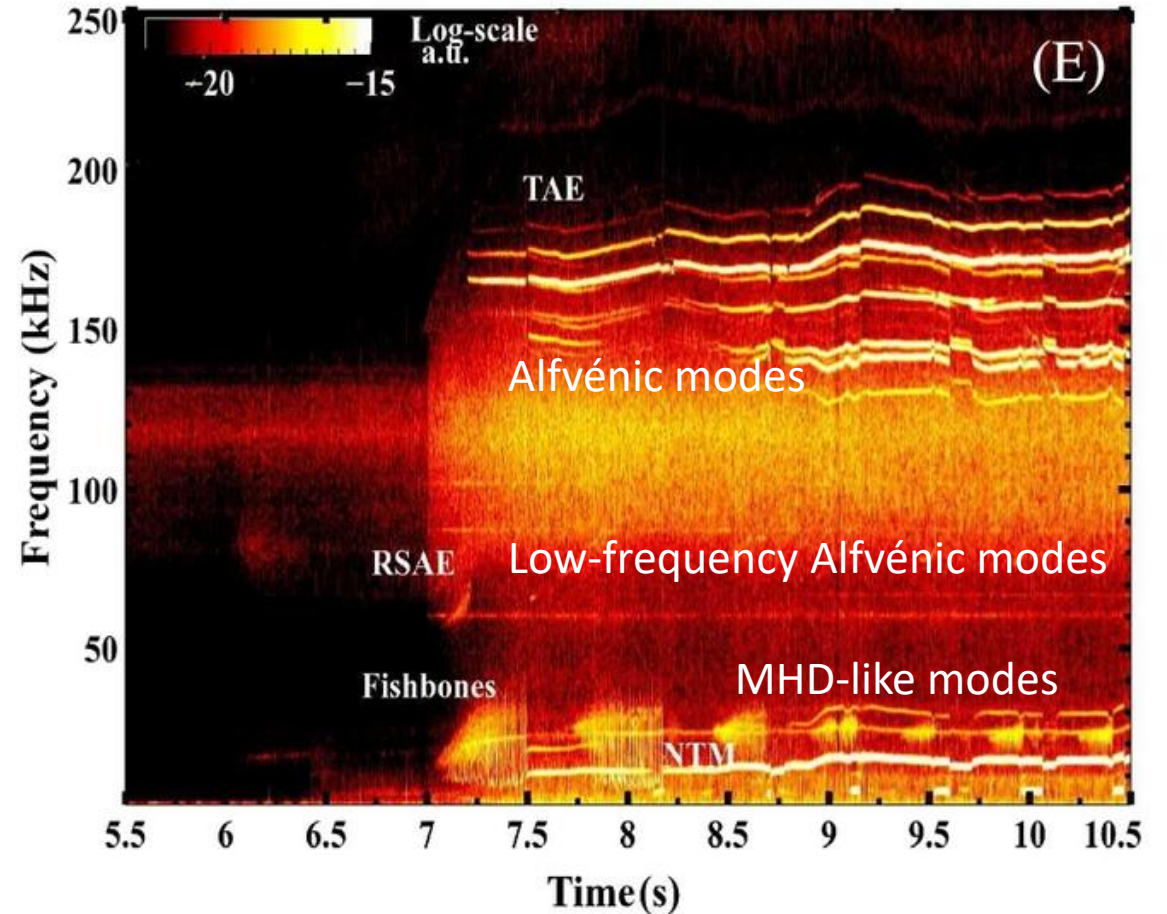
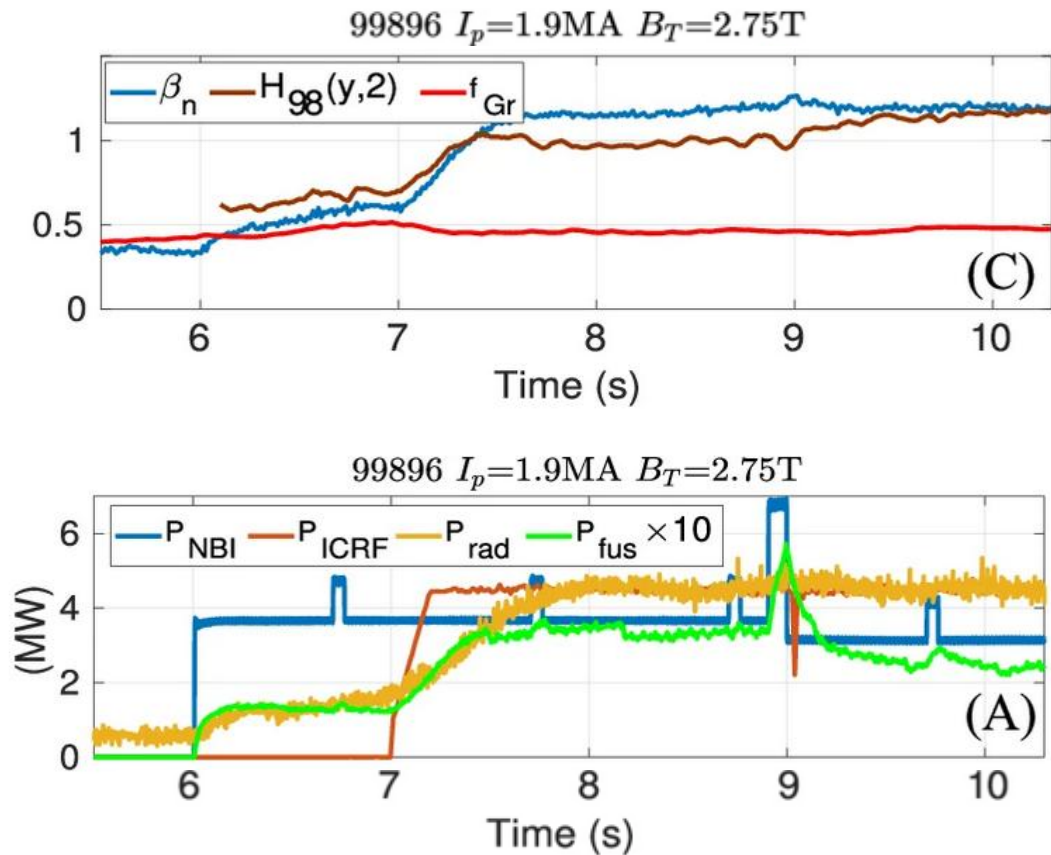


- Several experiments worldwide will produce burning plasmas (ITER, BEST, SPARC, Infinity, GIGA, Stellaris ...).
- **Energetic particles** will be important for cross-couplings between micro- and macro-scales.
- **Turbulence and zonal structures** cannot be considered separately from **fast particles and Alfvénic modes**.
- MHD-type instabilities will also affect the plasma dynamics, turbulence, and transport in phase space.
- Plasma profile evolution becomes a **multi-scale nonlinear problem with global and kinetic contributions**.

EUROfusion (TSVV-G) program:

- **Global nonlinear gyrokinetics** as a minimal inclusive description addressing this problem within a single framework.
- **Lower fidelity (reduced models):** hybrid particle-MHD codes, phase-space transport codes (quasilinear and PSZS).
- Developing theory and codes for future burning-plasma fusion devices and dedicated ongoing experiments.

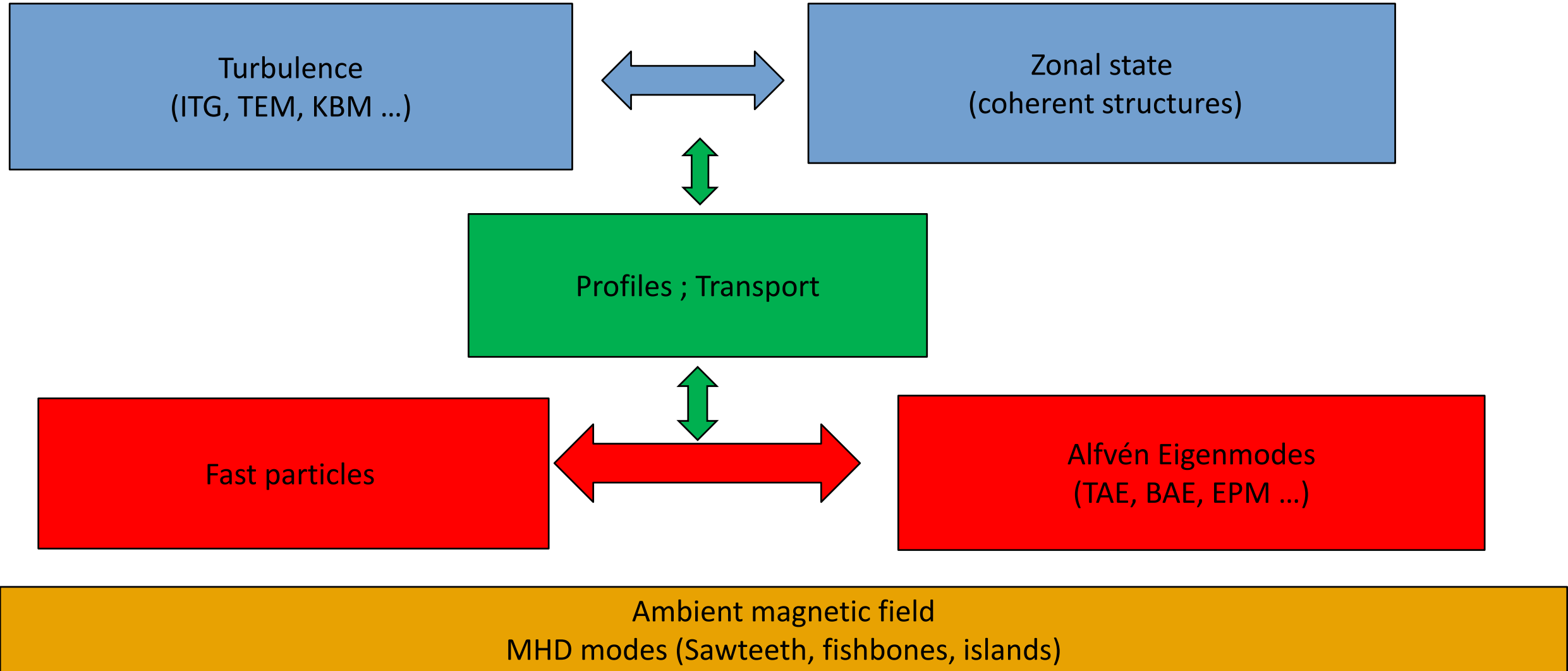
TSVV-G background: experiment (JET, Deuterium-Tritium campaign)



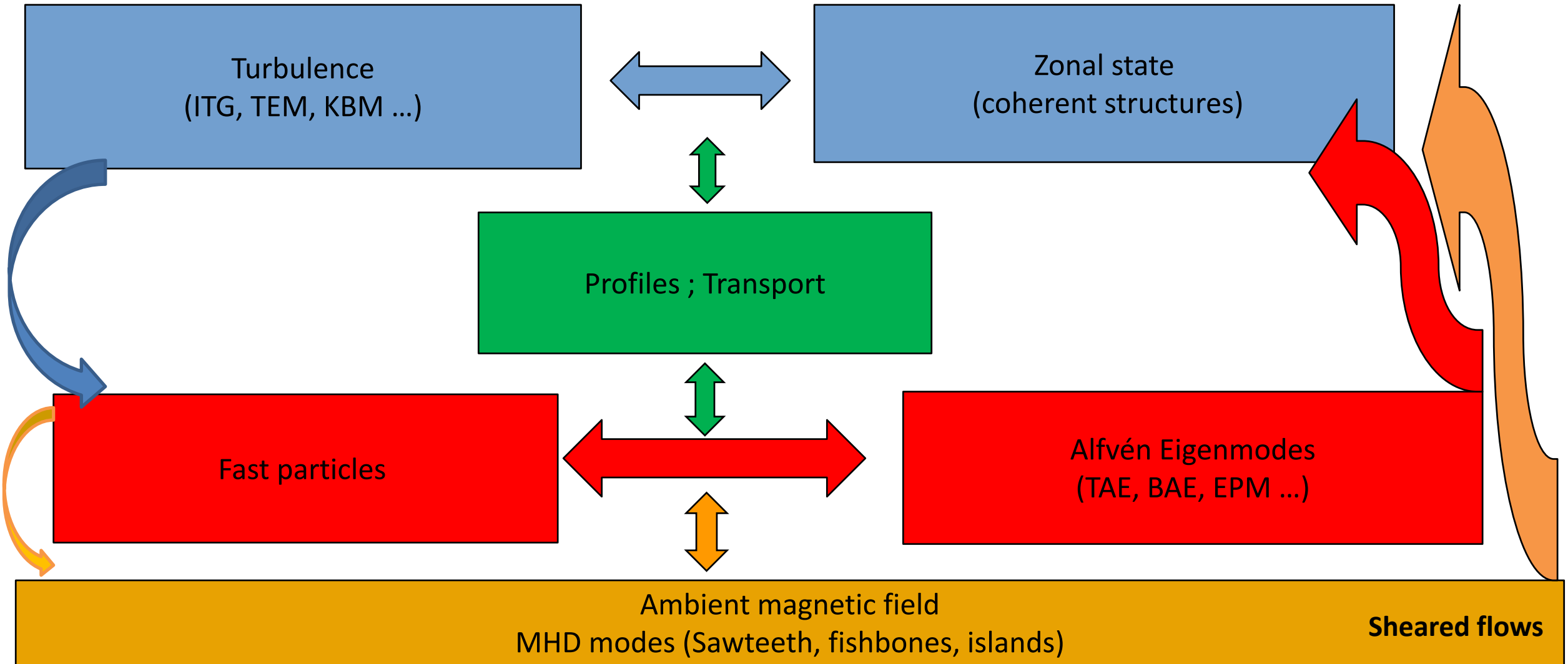
DT JET #99896: ICRF heating added at $t \sim 7\text{s}$; confinement improvement and increased fusion reaction (fast ions). Electromagnetic perturbations over a wide range of frequencies. Heat transport significantly reduced because of the energetic-particle driven instabilities.

[J. Garcia, Y. Kazakov, R. Coelho et al, Nature Communications \[2024\]](#)

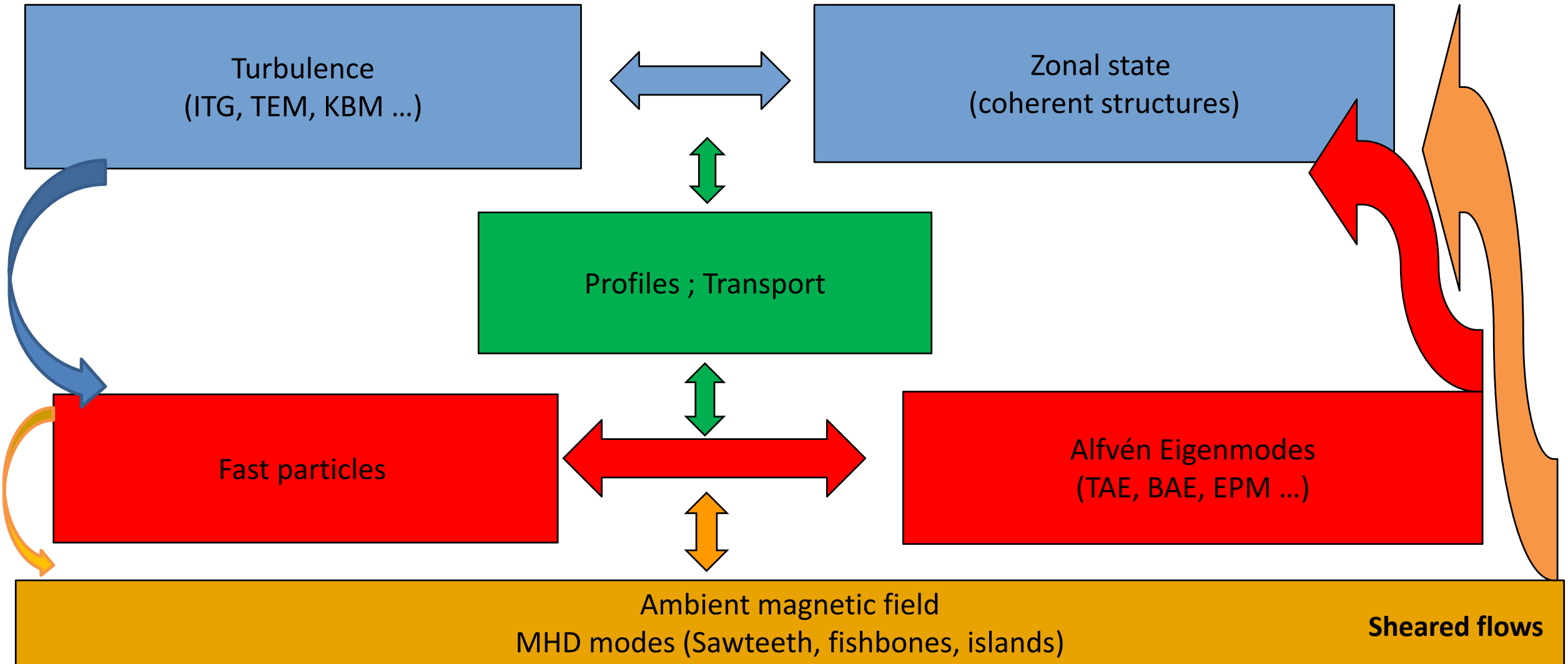
Synergies and couplings



Synergies and couplings



Burning plasmas: complex interconnected system



Global NL gyrokinetic theory is a minimal inclusive description

Gyrokinetic codes: ORB5 and EUTERPE codes



- ❖ Global nonlinear gyrokinetic particle-in-cell codes
- ❖ Electromagnetic, kinetic electrons, collision operators
- ❖ Arbitrary-wavelength solver, model for $\delta B_{||}$ available
- ❖ Both codes started at SPC EPFL Lausanne, closely related
- ❖ EUTERPE developed specifically for stellarator research
- ❖ Codes are supported by EUROfusion (e.g. E-TASC)
- ❖ Since June 2025: joint project, single repository
- ❖ The codes can further converge to a single codebase

Name	Last commit	Last update
euterpe	- updated supermuc modules i...	4 days ago
orb5	New TOK config	6 days ago

ORB5-EUTERPE community (EU)



Developers and users in many EU fusion sites
Contributions from India (IPR) and Ukraine

Crucial role played by EUROfusion!
(ENR projects 2014-2024, TSVV-10)



Cancellation problem

„Symplectic“ formulation of gyrokinetics:

$$\Gamma = q\mathbf{A}^* \cdot d\mathbf{R} + \frac{B}{\Omega} \mu d\theta + \langle A_{\parallel} \rangle \mathbf{b} \cdot d\mathbf{R} - \left(\frac{mv_{\parallel}^2}{2} + \mu B + q\langle \phi \rangle \right) dt$$

Problem: time derivative in equations of motion

$$\dot{v}_{\parallel} = -\frac{1}{m_s} \tilde{\mathbf{b}}^* \cdot \mu \nabla B - \frac{q_s}{m_s} \left(\tilde{\mathbf{b}}^* \cdot \nabla \langle \phi \rangle + \frac{\partial \langle A_{\parallel} \rangle}{\partial t} \right)$$

Not suitable for explicit solvers! Implicit too costly.

Solution: mixed-variable formulation:

$$A_{\parallel} = A_{\parallel}^{(s)} + A_{\parallel}^{(h)} \quad (\text{split the magnetic potential})$$

$$\Gamma = q\mathbf{A}^* \cdot d\mathbf{R} + \frac{m}{q} \mu d\theta + q \langle A_{\parallel}^{(s)} \rangle \mathbf{b} \cdot d\mathbf{R} - \left[\frac{mv_{\parallel}^2}{2} + \mu B + q \langle \phi - v_{\parallel} A_{\parallel}^{(h)} \rangle \right] dt$$

$$\frac{\partial}{\partial t} A_{\parallel}^{(s)} + \mathbf{b} \cdot \nabla \phi = 0 \quad (\text{use new „freedom degree“})$$

[\[Mishchenko et al, PoP 2017\]](#)

„Hamiltonian“ formulation of gyrokinetics:

$$\Gamma = q\mathbf{A}^* \cdot d\mathbf{R} + \frac{B}{\Omega} \mu d\theta - \left(\frac{mp_{\parallel}^2}{2} + \mu B + q \langle \phi - p_{\parallel} A_{\parallel} \rangle \right) dt$$

Problem: „fictitious“ terms in Ampere’s law

$$\sum_{s=i,e,f} \frac{\beta_s}{\rho_s^2} \langle \bar{A}_{\parallel} \rangle_s - \nabla_{\perp}^2 A_{\parallel} = \mu_0 \sum_{s=i,e,f} j_{\parallel 1s}$$

These terms result from „adiabatic distribution function“

$$\bar{H}_1 = q_s \langle \phi - v_{\parallel} A_{\parallel} \rangle, \quad F_e^{(\text{ad})} = F_{0e} e^{-\bar{H}_1/T_e} - F_{0e} \approx -\frac{q_e F_{0e}}{T_e} \langle \phi - v_{\parallel} A_{\parallel} \rangle$$

They correspond to the „adiabatic current“

$$\mu_0 \bar{j}_{\parallel s}^{(\text{ad})} = \mu_0 q_s \int v_{\parallel} F_s^{(\text{ad})} d^3v = \frac{\mu_0 n_{0s} q_s^2}{m_s} A_{\parallel} = \frac{\beta_s}{\rho_s^2} A_{\parallel}$$

These „fictitious“ terms must cancel (they do analytically)

Numerically: skin term is on grid; current represented by particles

$$\frac{\beta_e}{\rho_e^2} A_{\parallel} = \frac{\mu_0 n_0 e^2}{m_e} A_{\parallel}$$

Problem: „fictitious“ terms can be very large!





Gyrokinetic equations in mixed-variables formulation

The equations include the gyrokinetic Vlasov equation:

$$\frac{\partial f_{1s}}{\partial t} + \dot{\mathbf{R}} \cdot \frac{\partial f_{1s}}{\partial \mathbf{R}} + \dot{v}_{\parallel} \frac{\partial f_{1s}}{\partial v_{\parallel}} = - \dot{\mathbf{R}}^{(1)} \cdot \frac{\partial F_{0s}}{\partial \mathbf{R}} - \dot{v}_{\parallel}^{(1)} \frac{\partial F_{0s}}{\partial v_{\parallel}}, \quad (1)$$

the equations for the gyro-center orbits:

$$\dot{\mathbf{R}} = \dot{\mathbf{R}}^{(0)} + \dot{\mathbf{R}}^{(1)}, \quad \dot{v}_{\parallel} = \dot{v}_{\parallel}^{(0)} + \dot{v}_{\parallel}^{(1)} \quad (2)$$

$$\dot{\mathbf{R}}^{(0)} = v_{\parallel} \mathbf{b}_0^* + \frac{1}{qB_{\parallel}^*} \mathbf{b} \times \mu \nabla B, \quad \dot{v}_{\parallel}^{(0)} = - \frac{\mu}{m} \mathbf{b}_0^* \cdot \nabla B \quad (3)$$

$$\dot{\mathbf{R}}^{(1)} = \frac{\mathbf{b}}{B_{\parallel}^*} \times \nabla \langle \phi - v_{\parallel} A_{\parallel}^{(s)} - v_{\parallel} A_{\parallel}^{(h)} \rangle - \frac{q}{m} \langle A_{\parallel}^{(h)} \rangle \mathbf{b}_0^* \quad (4)$$

$$\dot{v}_{\parallel}^{(1)} = - \frac{q}{m} \left[\mathbf{b}^* \cdot \nabla \langle \phi - v_{\parallel} A_{\parallel}^{(h)} \rangle + \frac{\partial}{\partial t} \langle A_{\parallel}^{(s)} \rangle \right] - \frac{\mu}{m} \frac{\mathbf{b} \times \nabla B}{B_{\parallel}^*} \cdot \nabla \langle A_{\parallel}^{(s)} \rangle \quad (5)$$

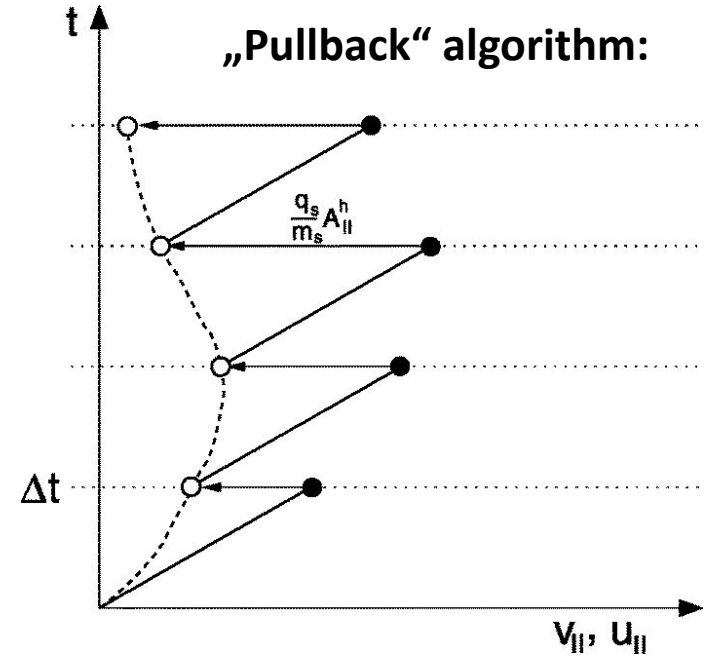
$$\mathbf{b}^* = \mathbf{b}_0^* + \frac{\nabla \langle A_{\parallel}^{(s)} \rangle \times \mathbf{b}}{B_{\parallel}^*}, \quad \mathbf{b}_0^* = \mathbf{b} + \frac{mv_{\parallel}}{qB_{\parallel}^*} \nabla \times \mathbf{b} \quad (6)$$

$$B_{\parallel}^* = B + \frac{mv_{\parallel}}{q} \mathbf{b} \cdot \nabla \times \mathbf{b}, \quad (7)$$

Field equations:

$$\frac{\partial}{\partial t} A_{\parallel}^{(s)} + \mathbf{b} \cdot \nabla \phi = 0, \quad - \nabla \cdot \left(\frac{n_0}{B\omega_{ci}} \nabla_{\perp} \phi \right) = \bar{n}_{1i} - \bar{n}_{1e}$$

$$\sum_{s=i,e} \frac{\beta_s}{\rho_s^2} A_{\parallel}^{(h)} - \nabla_{\perp}^2 A_{\parallel}^{(h)} = \mu_0 \sum_{s=i,e} \bar{j}_{\parallel 1s} + \nabla_{\perp}^2 A_{\parallel}^{(s)}$$



At the end of every time step:

1) re-split magnetic potential

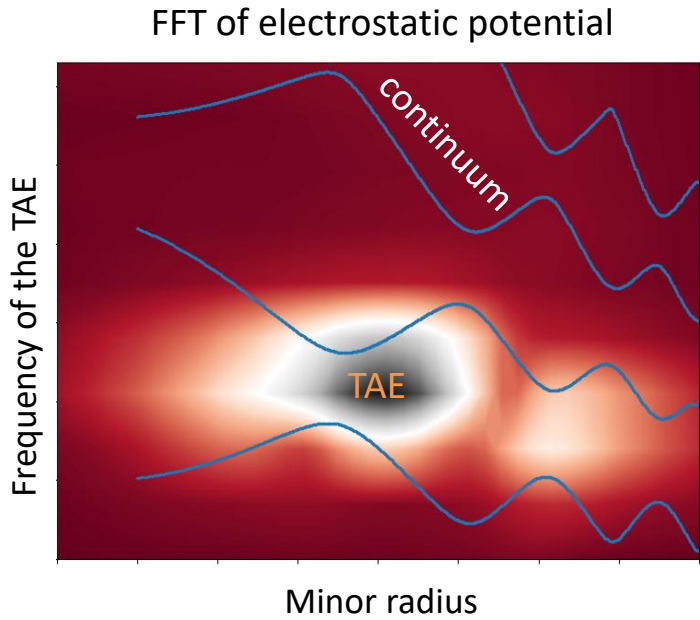
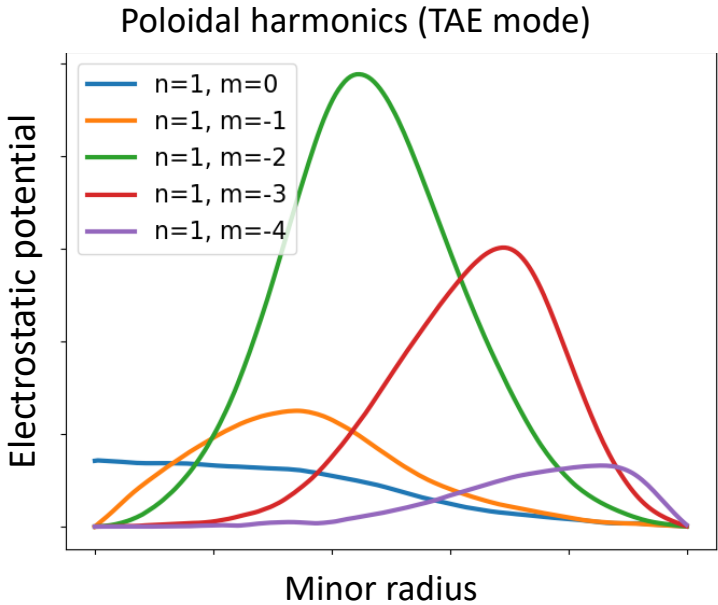
$$A_{\parallel(\text{new})}^{(s)}(t_i) = A_{\parallel}(t_i) = A_{\parallel(\text{old})}^{(s)}(t_i) + A_{\parallel(\text{old})}^{(h)}(t_i)$$

$$A_{\parallel(\text{new})}^{(h)}(t_i) = 0$$

2) re-define markers

$$f_1^{(s)}(Z^{(s)}, A_{\parallel}) = f_1^{(m)}(Z^{(m)}, A_{\parallel}^{(s)}, A_{\parallel}^{(h)})$$

$$v_{\parallel}^{(s)} = v_{\parallel}^{(m)} - \frac{e}{m} \langle A_{\parallel}^{(h)} \rangle$$



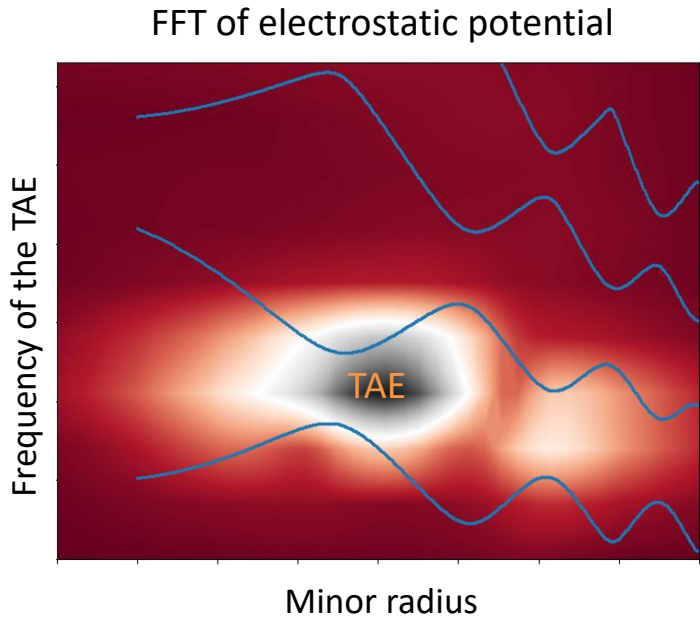
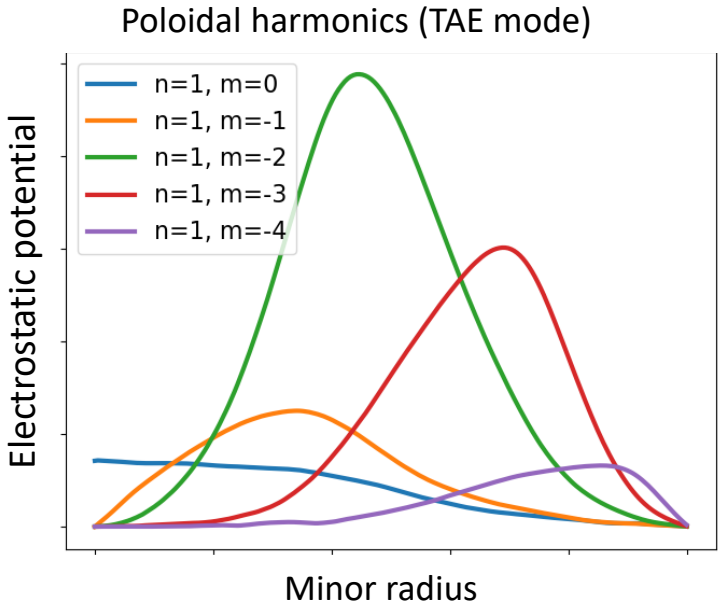
Energetic-particle driven TAE/EPM instabilities in TCV-shaped tokamak.

Coupled poloidal harmonics $m = -2$ and $m = -3$. Eigenmode frequency in the toroidal gap.

Gyrokinetic simulations of chirping instabilities: wave-particle nonlinearity



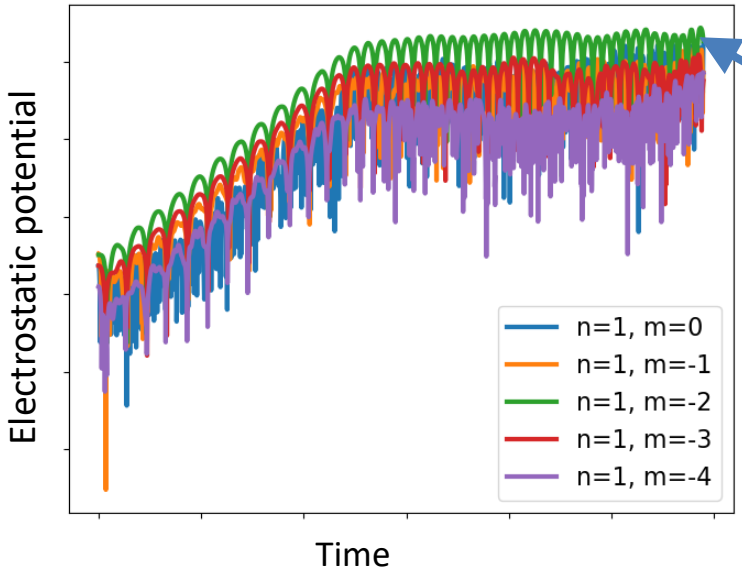
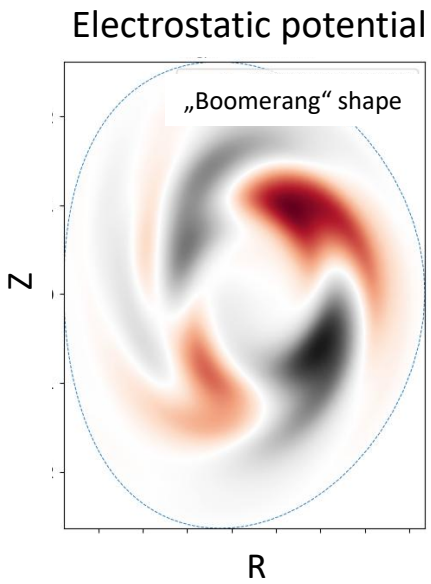
ORB5



Energetic-particle driven TAE/EPM instabilities in TCV-shaped tokamak.

Coupled poloidal harmonics $m = -2$ and $m = -3$. Eigenmode frequency in the toroidal gap.

“Boomerang” shape of 2D mode structure caused by energetic-particle resonant response.

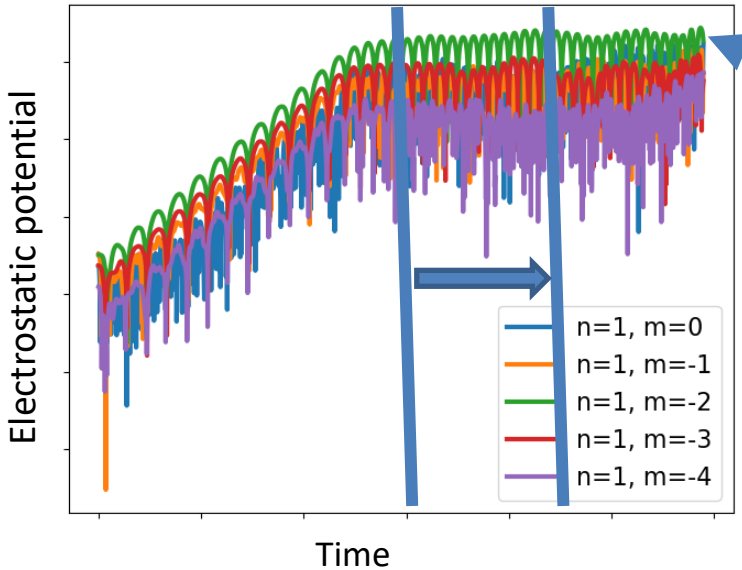
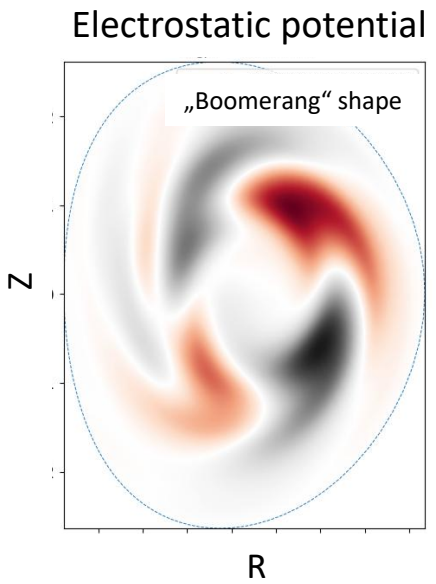
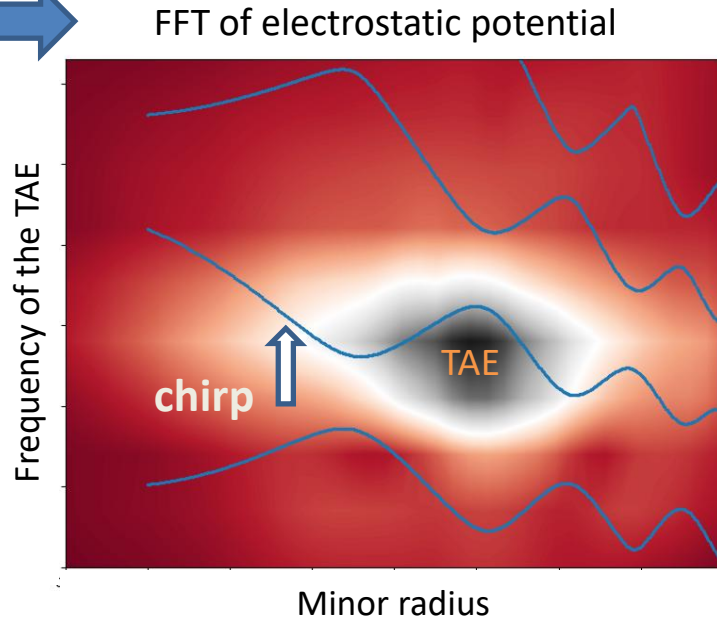
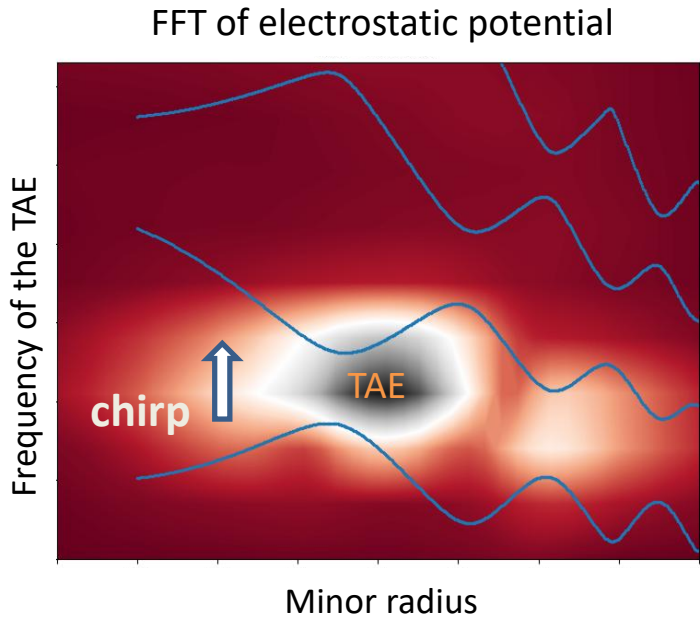
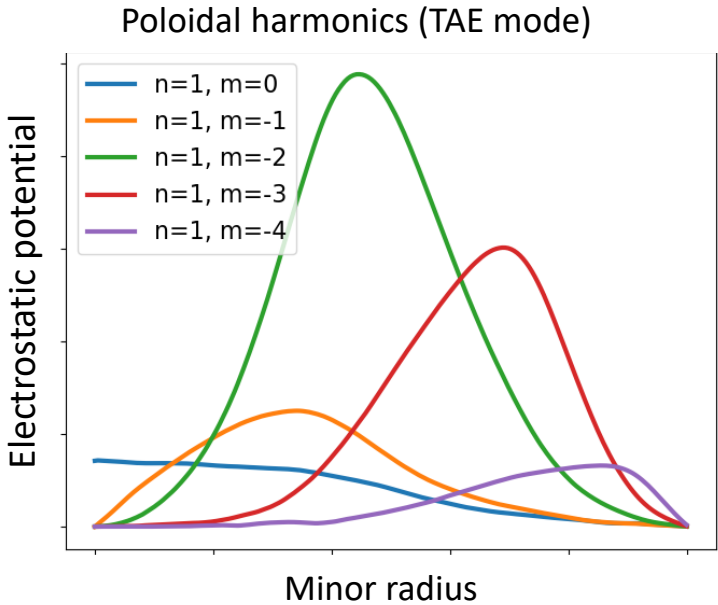


Saturation amplitude of the instability determines energetic-particle transport.

Gyrokinetic simulations of chirping instabilities: wave-particle nonlinearity



ORB5



Saturation amplitude of the instability determines energetic-particle transport.

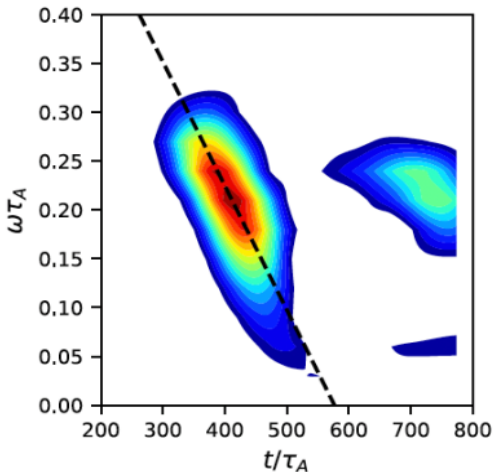
Chirping dynamics is observed for mode frequency: phase-space structures.

Gyrokinetic simulations of chirping instabilities: wave-particle nonlinearity



ORB5

Comprehensive studies of the nonlinear frequency on Alfvén modes driven by fast ions. A large number of global gyrokinetic simulations performed with ORB5. Plasma parameters, energetic-particle parameters varied. Chirping rate follows universal dependency.



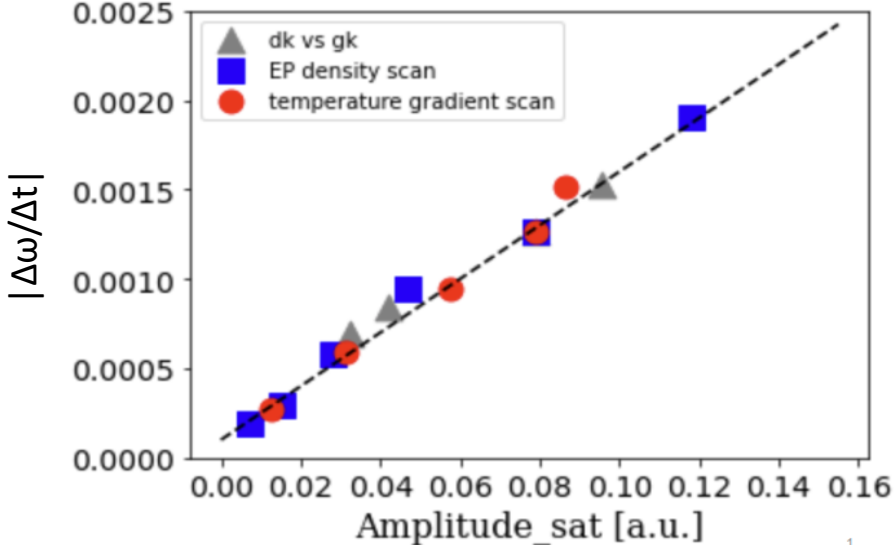
Time evolution of the mode frequency

various parameter scans to map chirping rate vs saturation amplitude



Energetic particle density, energy, ...

$$\Delta\omega/\Delta t \sim A_{\text{sat}}$$

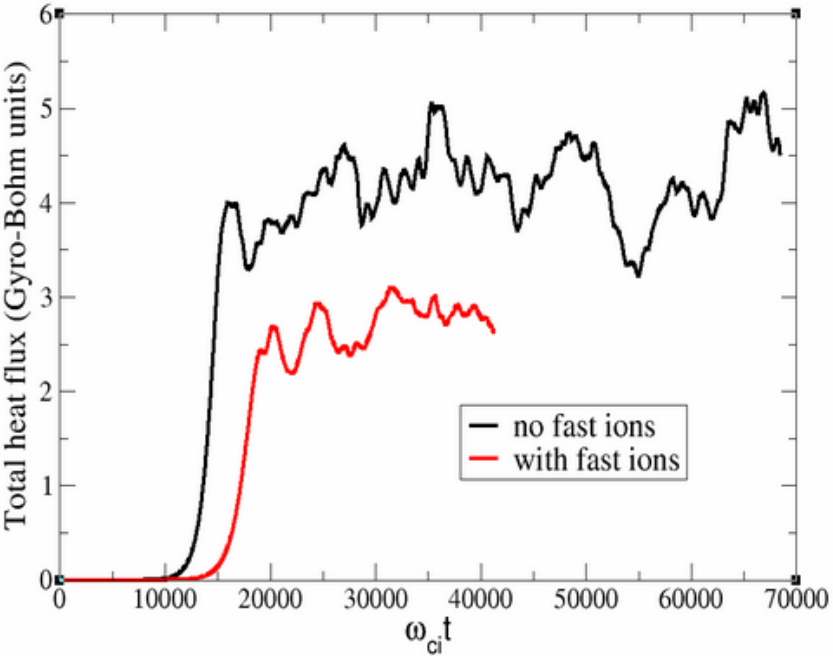
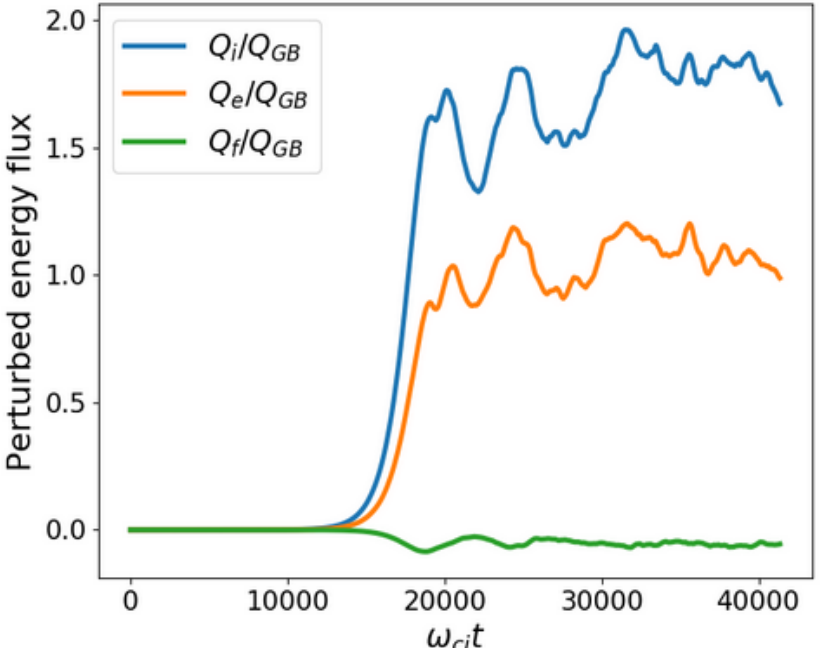


Theory compared to simulations (verification): chirping modes in presence of energetic particles. [Chirping rate from simulations is proportional to saturation amplitude \(as predicted theoretically, \[Zonca et al., Varenna Conf. 2024\]\)](#)

ORB5 code: effect of fast ions on core plasma energy confinement



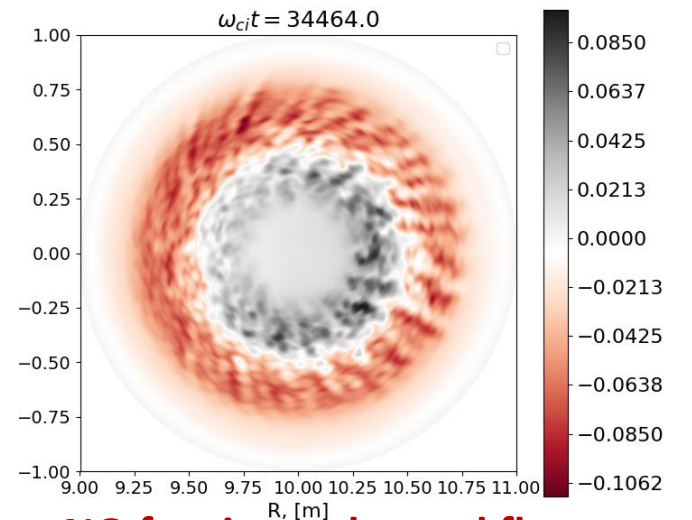
ORB5



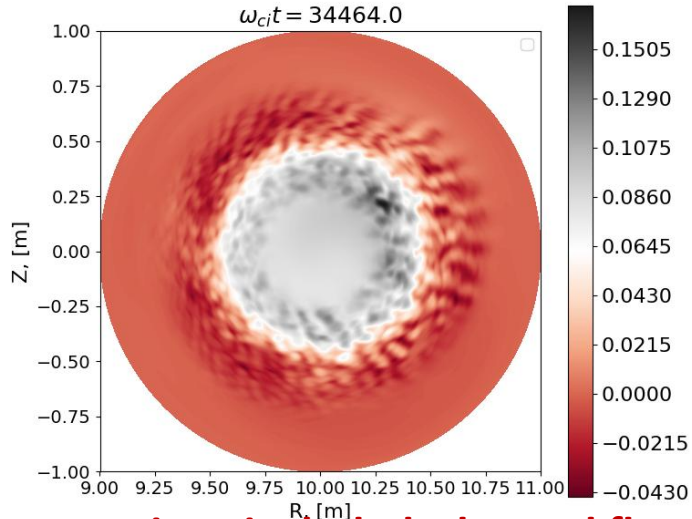
For $\beta_e = 0.1\%$, fast-ion heat flux is small.

Total heat flux reduced by the fast ions!

Nonlinear gyrokinetic simulations with ORB5 code demonstrate clear fast-particle reduction in the heat flux for both the bulk ions and the electrons
(in agreement with earlier flux-tube results) $\beta_e = 0.1\%$



NO fast ions: sheared flow



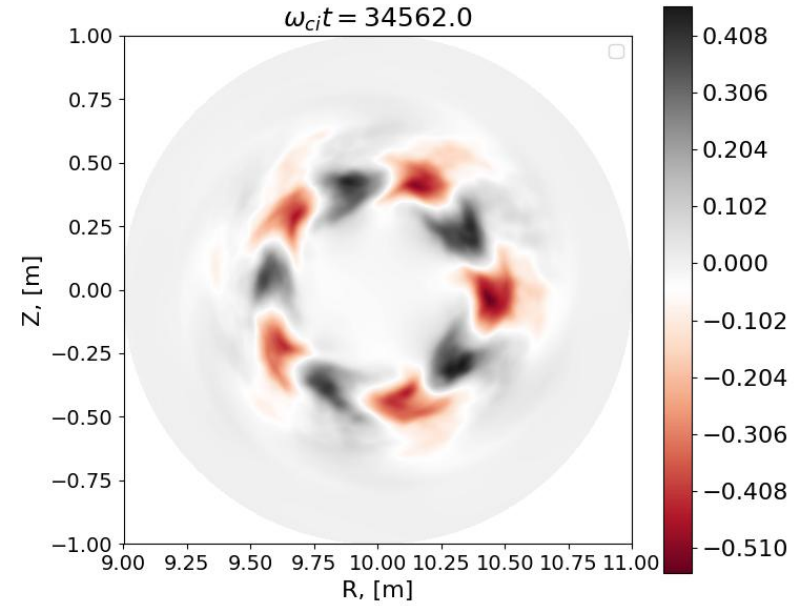
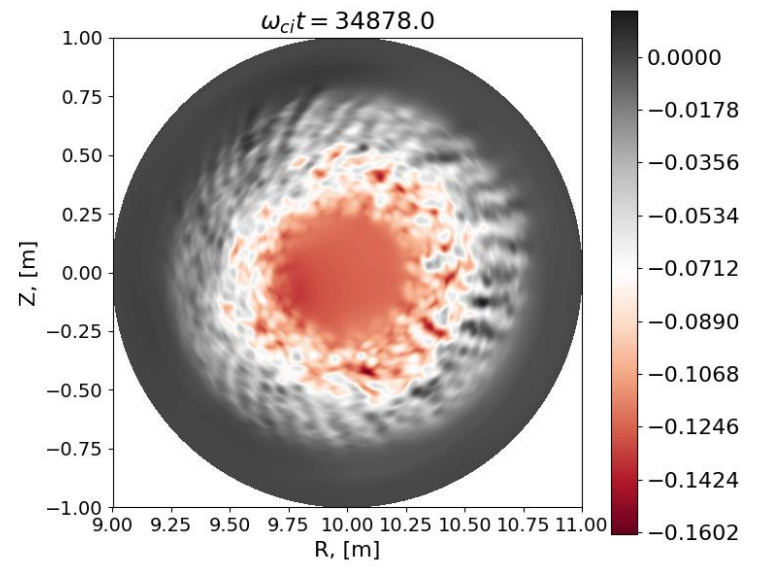
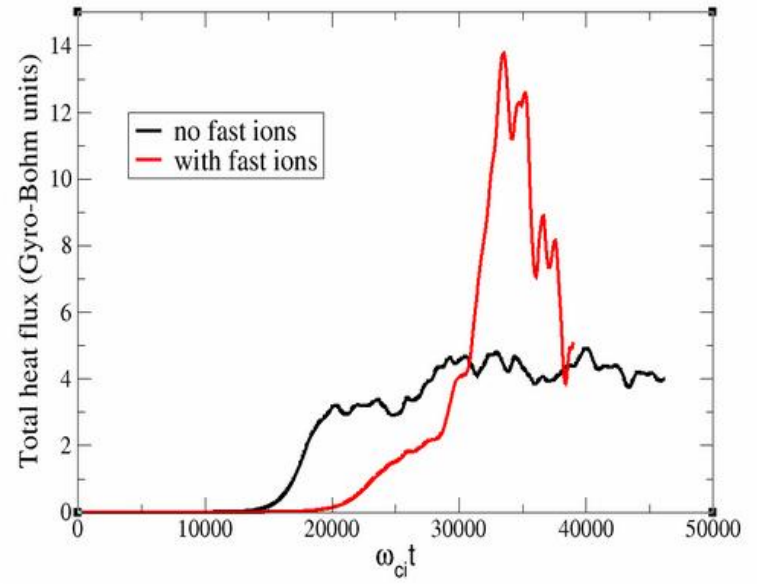
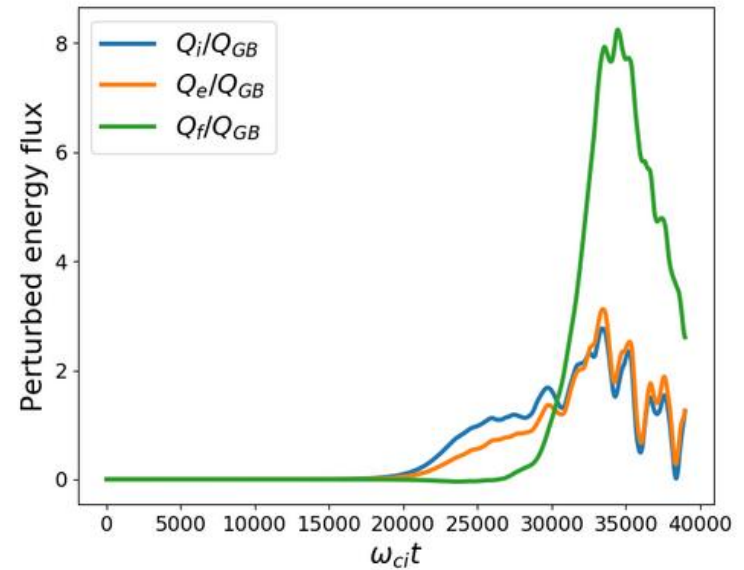
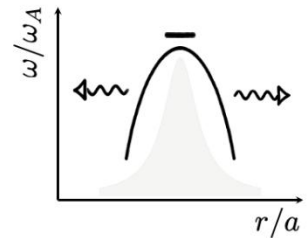
Fast ions included: sheared flow



ORB5 code: effect of fast ions on core plasma confinement

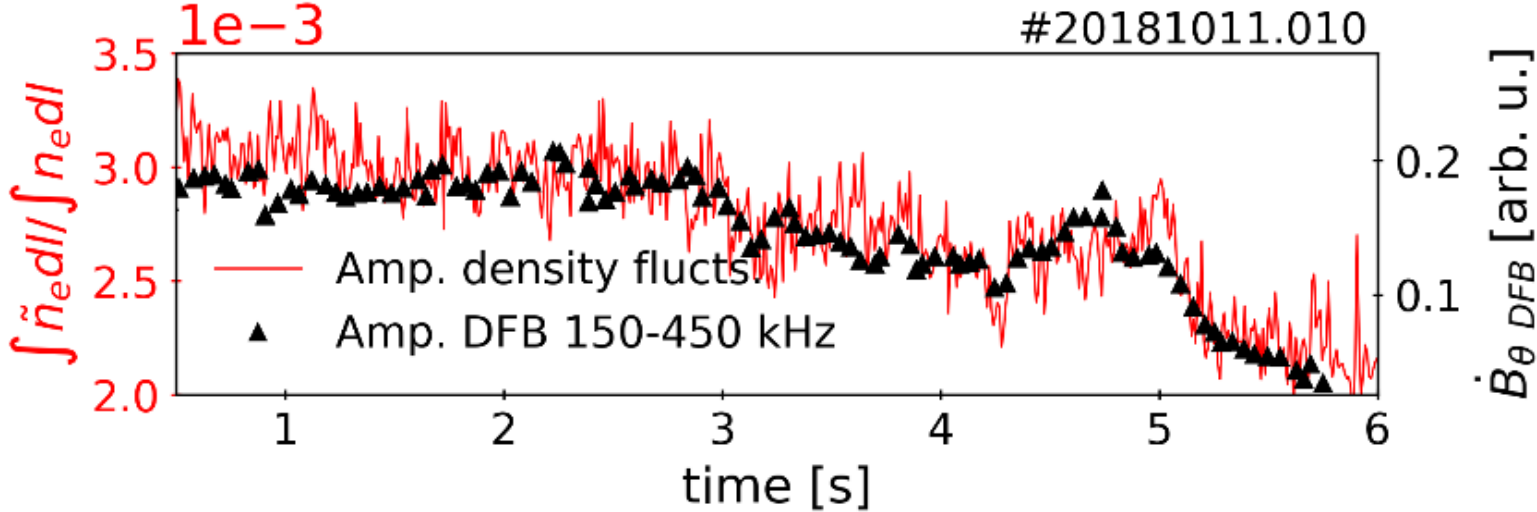
- For $\beta_e = 0.24\%$, the dynamics is different. Fast ion heat flux is substantial. Total heat flux is not reduced
- Global Alfvénic mode (a BAE) develops driving fast ion energy flux. Work in progress!
- Properties of the Zonal State become increasingly more global as plasma β is increased
- Mesoscale phenomena, such as kinetic Alfvén wave propagation, need to be captured for a correct description of the plasma dynamics

ORB5



No fast ions: sheared flow

Fast ions included: global (BAE) mode

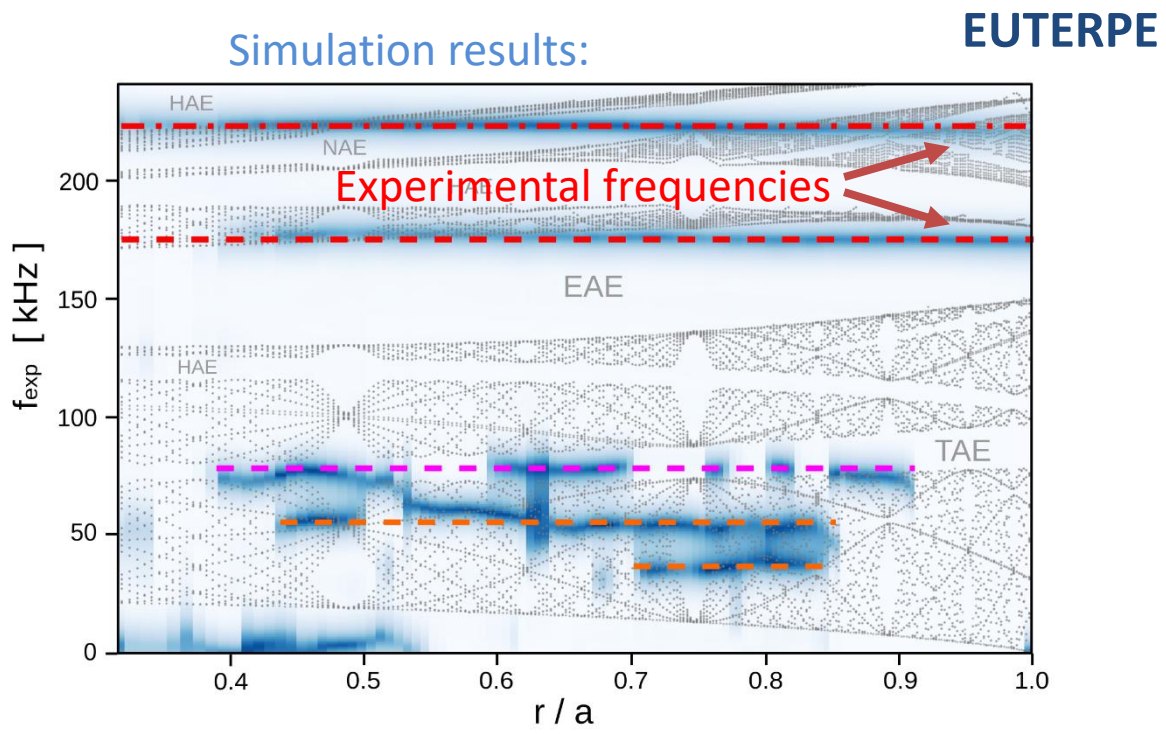
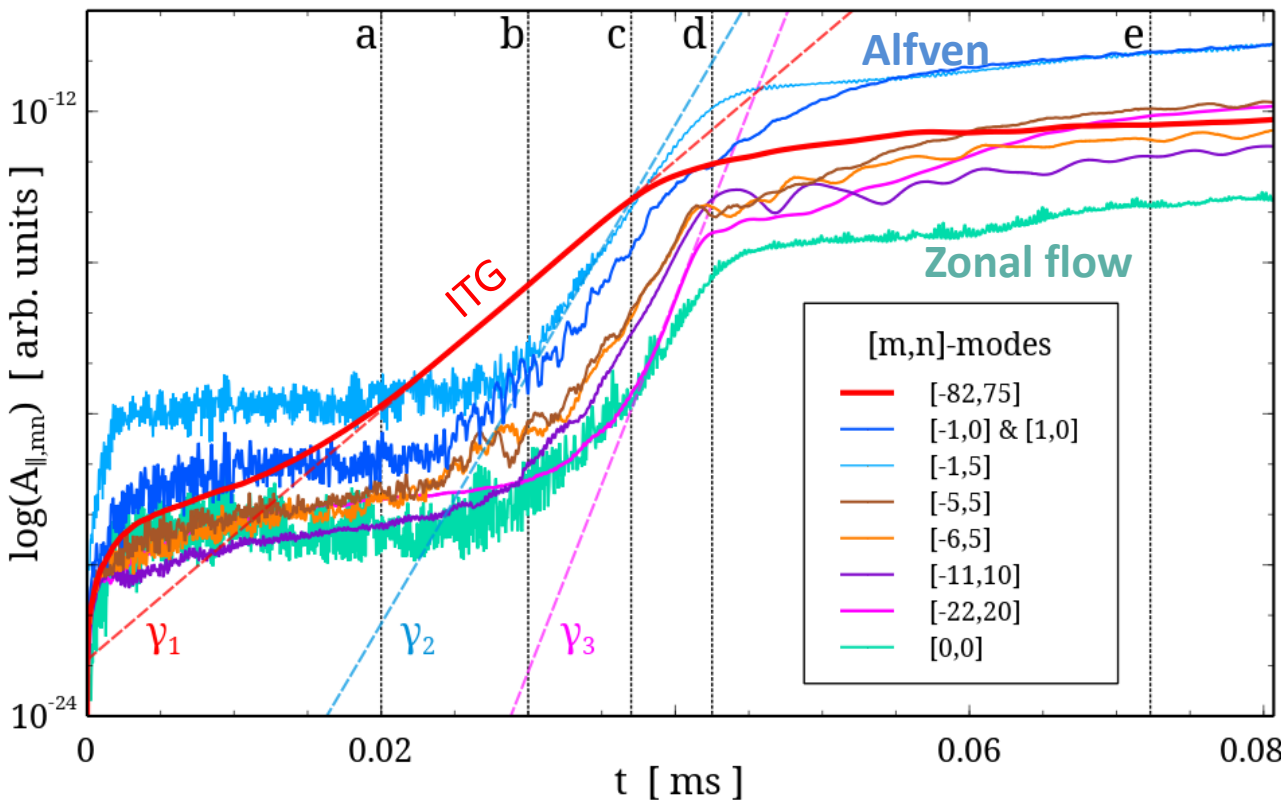


Experimental Results:

- predominant ITG turbulence observed in W7-X
- distinctive Alfvénic frequency band is often seen

- temporal evolution of turbulence (PCI) and magnetic fluctuations (Mirnov) is very similar over entire discharges
- Alfvénic fluctuation amplitudes and turbulence level are correlated for many plasma discharges

Hypothesis: Alfvénic fluctuations are nonlinearly driven by ITG turbulence in W7-X.



- early ITG (high-m) activity
- Zonal Flow & Alfvénic modes when ITG amplitude has passed threshold
- growth rate cascade $\gamma_1 : \gamma_2 : \gamma_3 = \gamma : 2\gamma : 3\gamma$ (similar to “beat-driven” ZFs)

Gyrokinetic simulations recover experimental frequencies

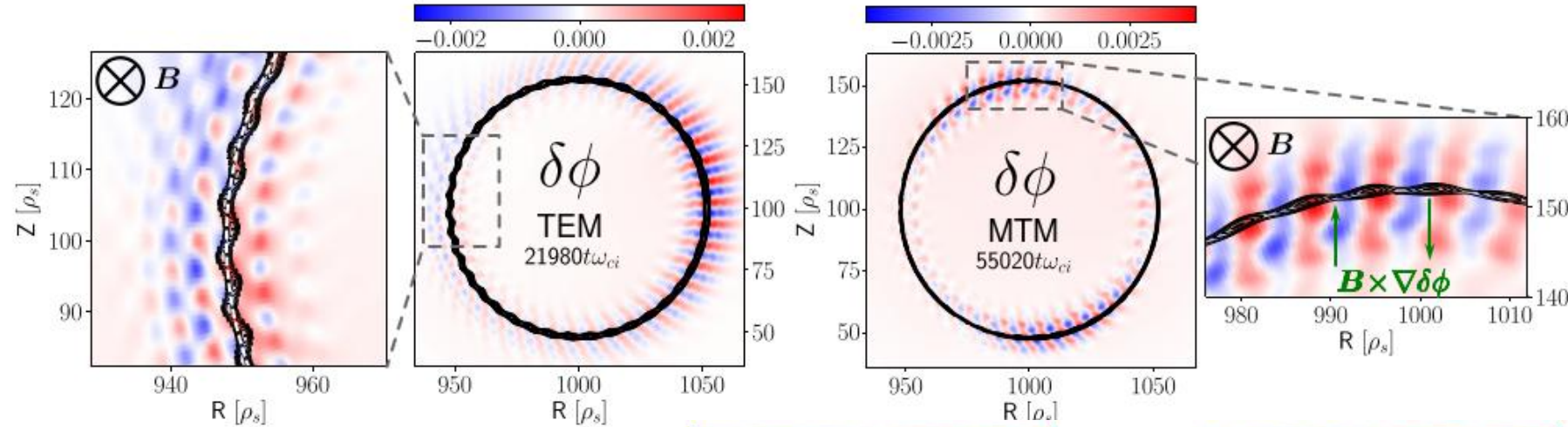
The role of ITG activity:

1. ITG modes included in simulations → low-m Alfvénic modes nonlinearly excited
2. ITG modes excluded from simulations → low-m modes remain stable



ORB5

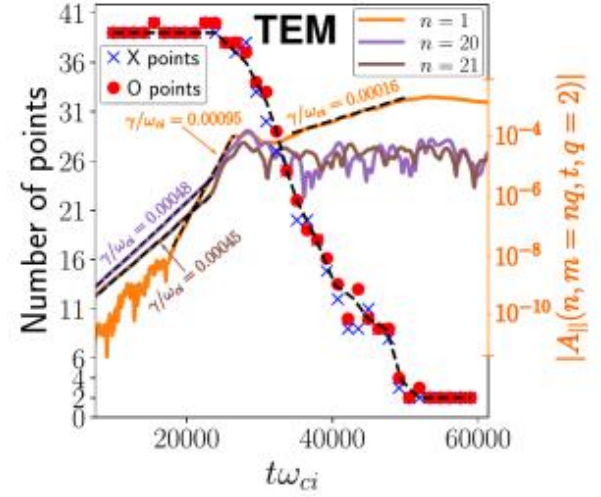
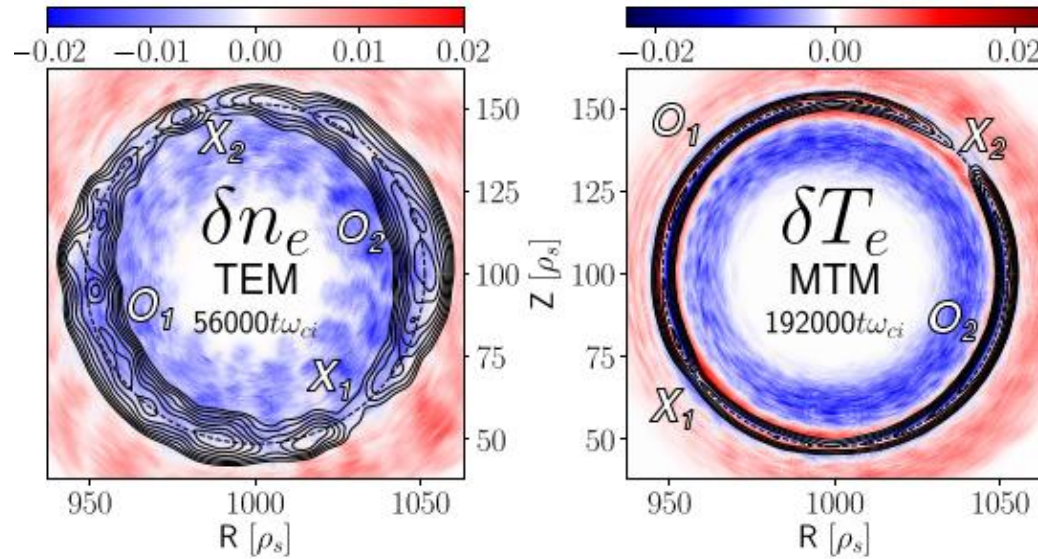
Turbulence-Driven Magnetic Islands in Toroidal Geometry



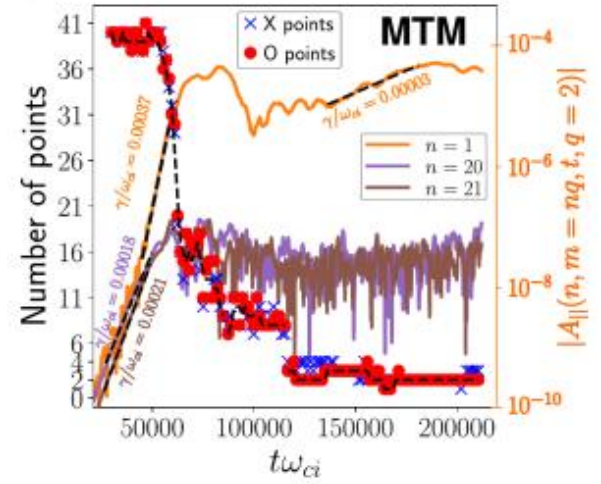
Linearly stable tearing mode in toroidal geometry (collisionless).

Micro-instabilities generate an $E \times B$ flow that drives magnetic field line reconnection forming multiple small-scale islands along resonant surface.

The small-scale islands interact non-linearly and eventually coalesce into a large-scale magnetic island.



(a) TEM case, $-R/L_n = 5$, $-R/L_{Ts} = 0$.

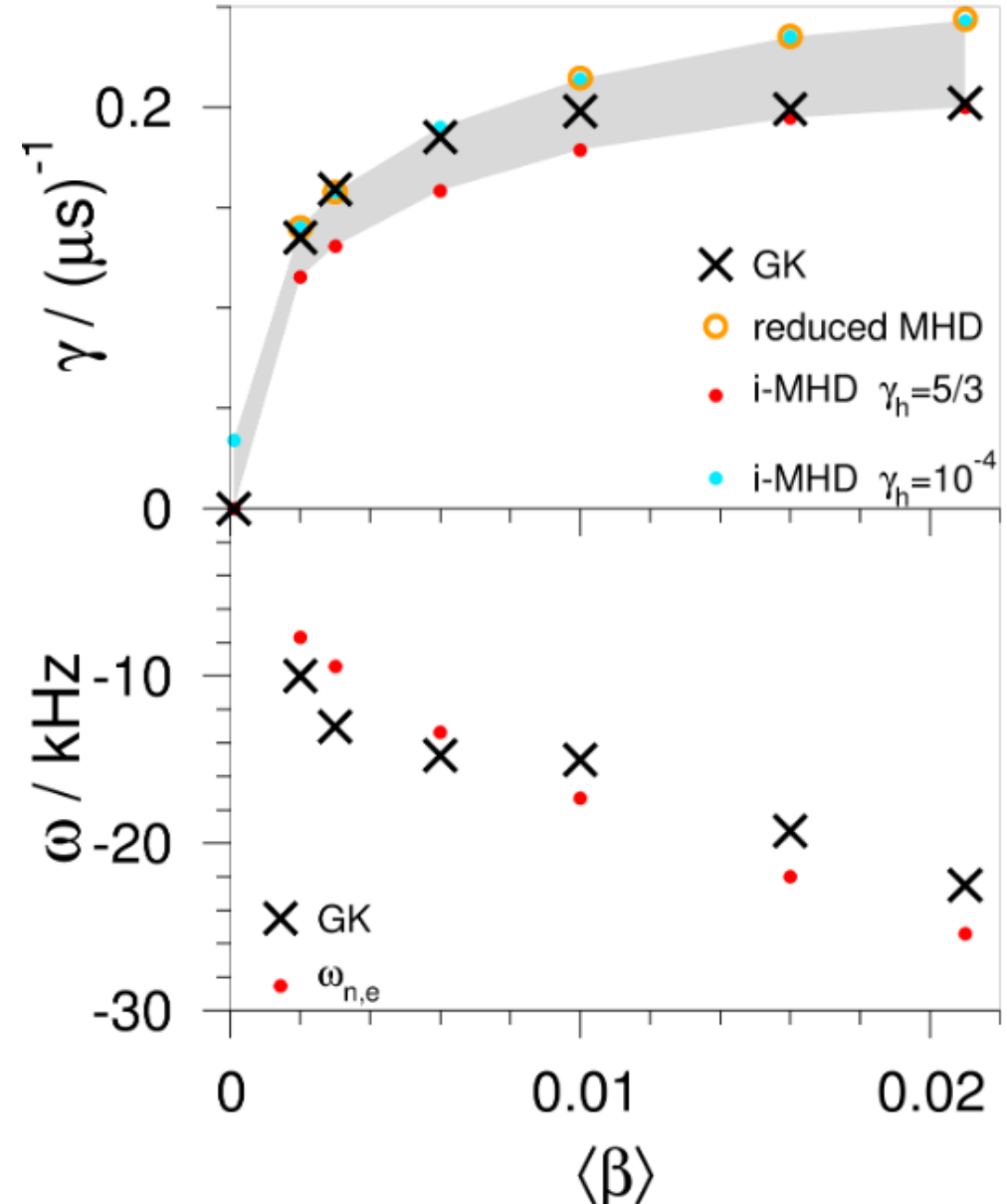




Gyrokinetic simulations of MHD-unstable stellarator configurations

- Consider 4-period stellarators: turning-ellipse configurations
- Choice: Mercier-unstable at finite β (not a reactor candidate)
- Perform full-torus GK simulations
- Find good agreement with MHD (CAS3D, CKA) in linear phase.

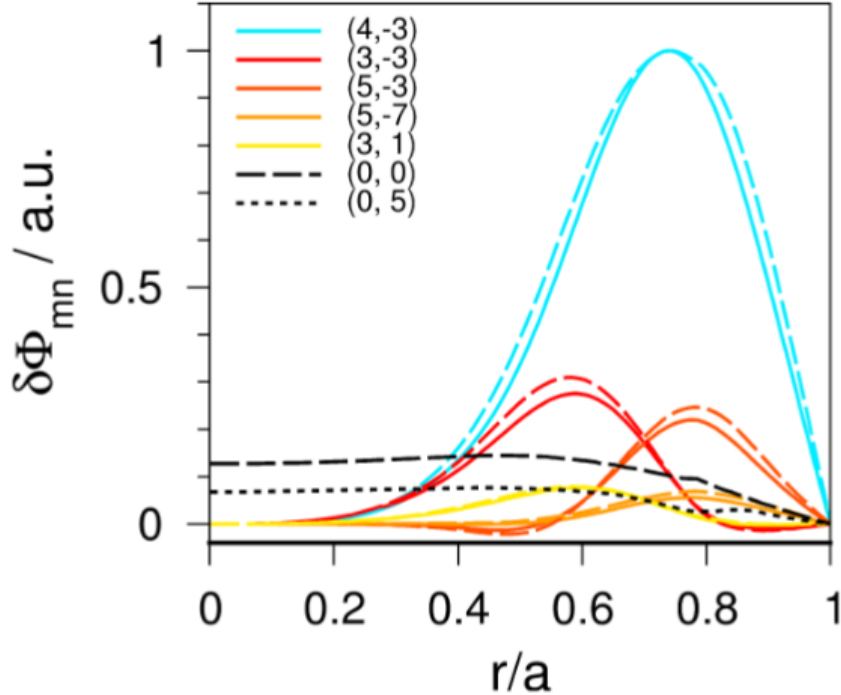
First-ever benchmark of global MHD with global gyrokinetics in stellarators!



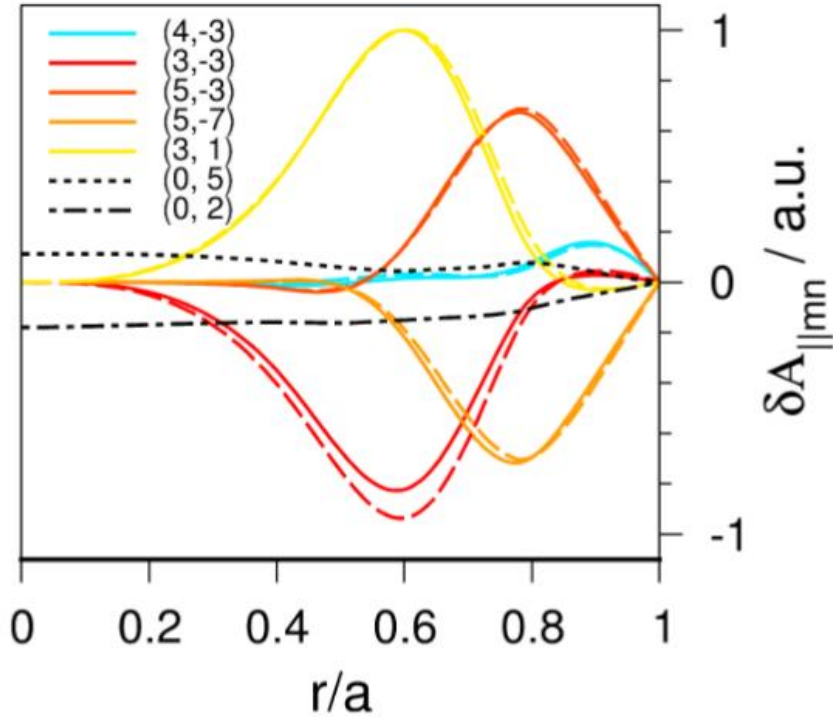


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First-ever benchmark of global MHD with global gyrokinetics in stellarators!



The ideal-MHD results (solid lines, $\gamma_h = 5/3$), are compared to the gyrokinetic results (dashed).



Very good agreement!



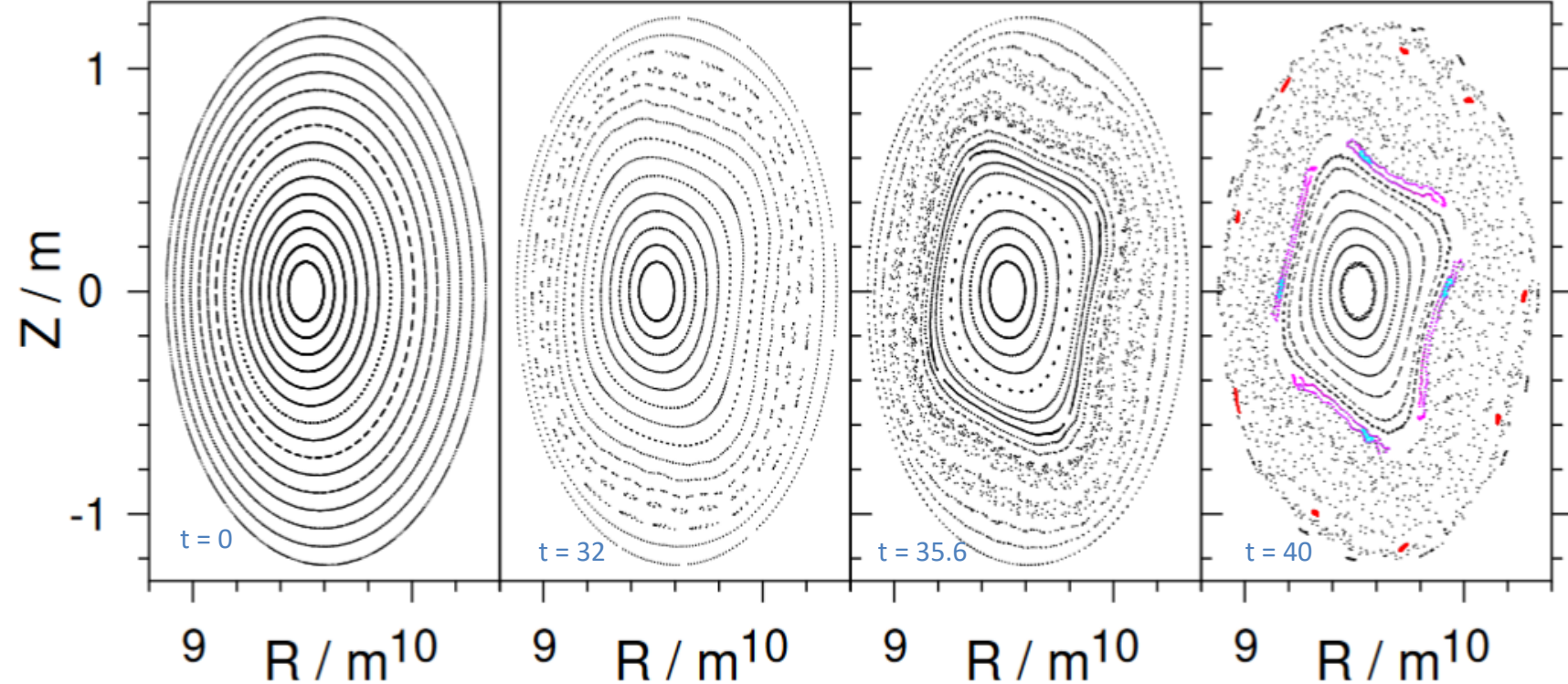
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First-ever benchmark of global MHD with global gyrokinetics in stellarators!

**Soft decay of magnetic surfaces.
No disruptive dynamics!
Soft beta limit in stellarators.**

EUTERPE

Perturbed field-line tracing for GK magnetic field:
nonlinear growth of magnetic islands, ergodization
 $\langle \beta \rangle = 0.006$



MHD instabilities in stellarator destroy flux surfaces.
Magnetic configuration is affected.

First-ever fully gyrokinetic simulations of nonlinear MHD in stellarator plasmas!

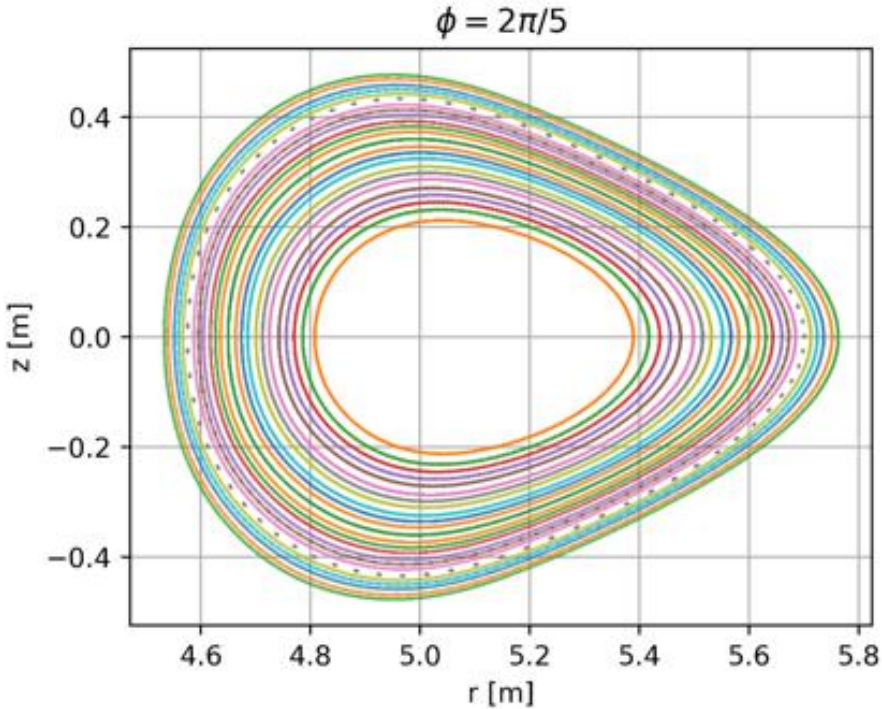
Gyrokinetic simulations of turbulence in W7-X: magnetic configuration affected



EM turbulence in Mercier-unstable W7-X configuration at $\beta = 4\%$.

EUTERPE

Turbulent evolution of total magnetic field (flux surfaces).

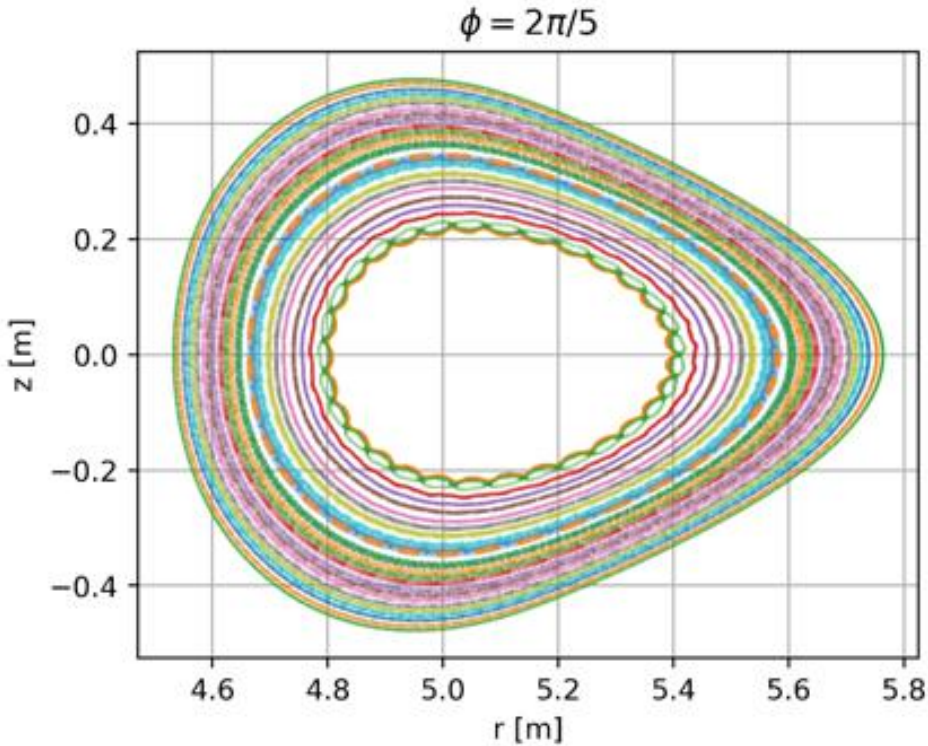


Linear phase

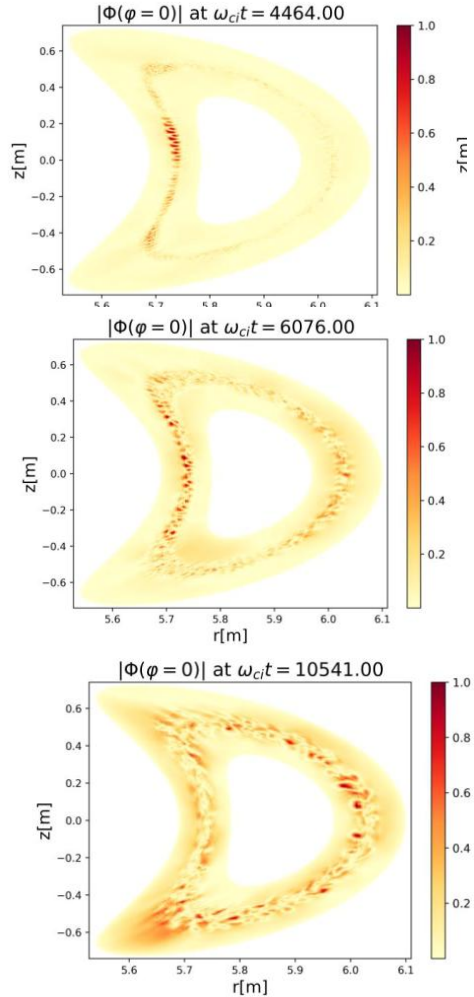


Turbulence leads to ergodization of flux surface.

Magnetic configuration is affected by EM turbulence.



Nonlinear phase



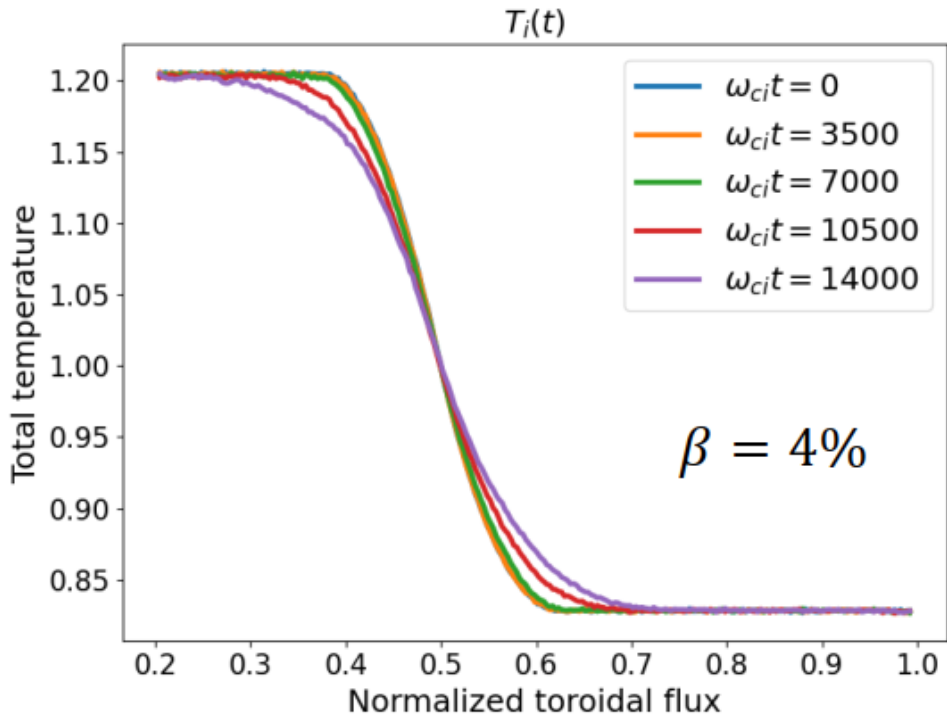
Gyrokinetic simulations of turbulence in W7-X: magnetic configuration affected



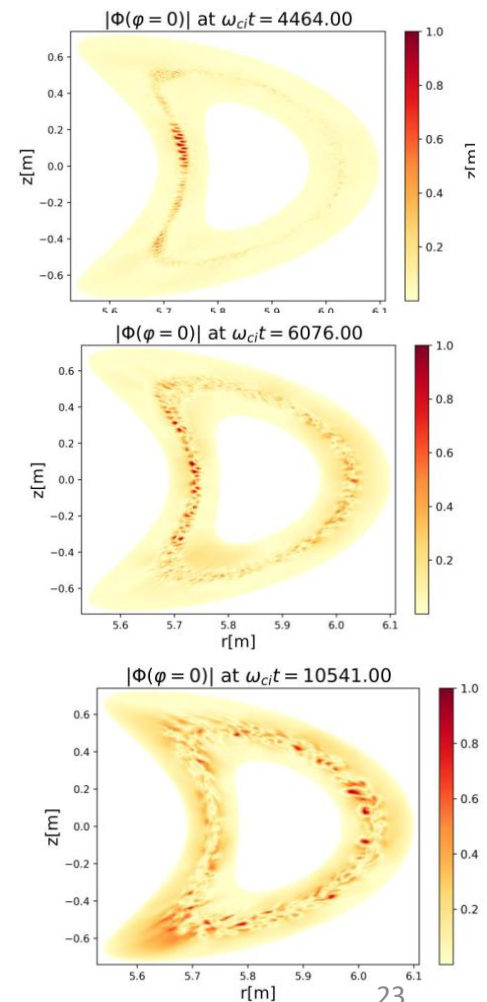
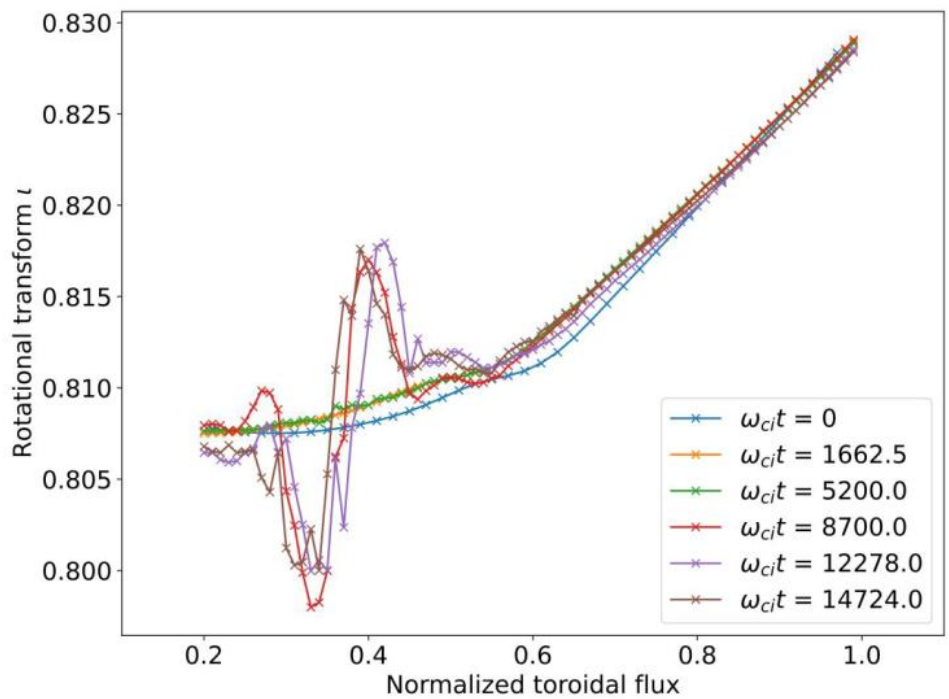
EM turbulence in Mercier-unstable W7-X configuration at $\beta = 4\%$.

EUTERPE

Turbulent evolution of temperature and safety factor profiles



Turbulence modifies rotation transform (safety factor).
Magnetic configuration is affected by EM turbulence.





- Turbulence, zonal flows, Alfvénic modes, energetic particles, MHD-like perturbations coexist and affect each other
- Global nonlinear gyrokinetics is a minimal inclusive approach for this system
- ORB5 and EUTERPE codes provide such a framework; enabled by numerous improvements in algorithms
- First-in-kind simulations of nonlinear plasma evolution performed in tokamak and stellarator geometries
- Hybrid models (kinetic-fluid) and reduced models (quasilinear, PSZS) have also been developed
- EUROfusion software ecosystem will make a qualitative difference for future fusion reactor projects (public and private)

<https://indico.euro-fusion.org/category/283/>

Outlook (TSVV-G):

New contributors: GYSELA and GENE in addition to ORB5/EUTERPE – all major GK codes represented; cooperation between teams, exchange of experience, favourable environment for scientific discovery

Dissemination: small associations (Greece and Ukraine) included in TSVV-G

User communities for major codes (beyond TSVVs): XTOR-K, HYMAGYC, ORB5/EUTERPE, ...



Eulerian code GENE and semi-Lagrangian full-f code GYSELA join PIC codes ORB5/EUTERPE: verification and leveraging each others strengths as elements of the overall EUROfusion (E-TASC) ecosystem

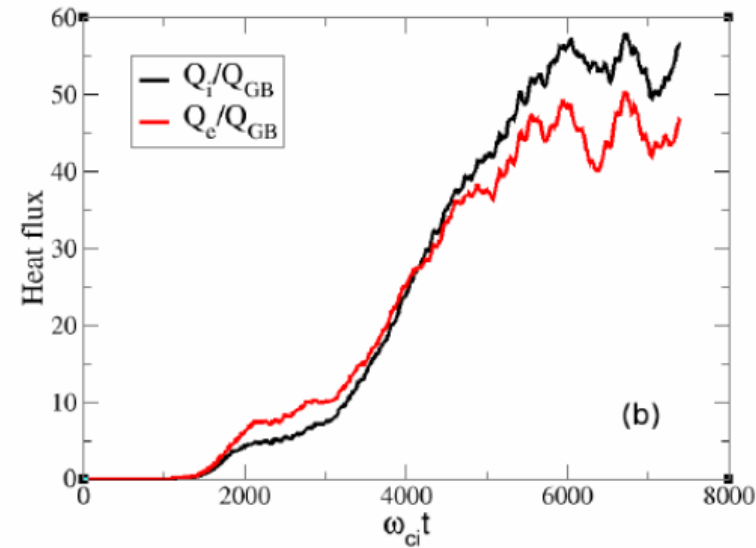
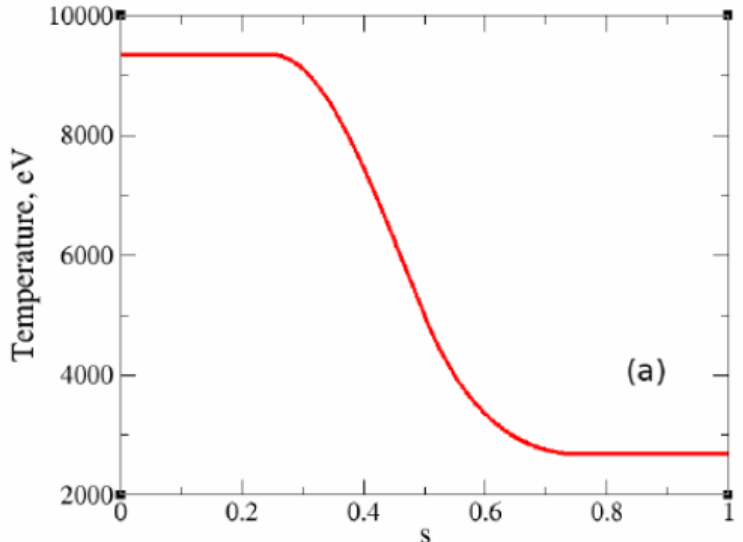
ORB5/EUTERPE applications:

- Saturation amplitude studies using ORB5: AEs + EPs (+ turbulence)
- Effect of EPs/AEs on turbulence (e.g. via ZF generation): ORB5
- Kink/fishbone instability; effect of ZFs, interaction with turbulence: ORB5 and GENE
- Tearing instability and turbulence in tokamaks (turbulent cascades, island merging vs. island healing), in realistic geometry and including collisions: ORB5 and GENE
- Electromagnetic turbulence in stellarators, EP physics: EUTERPE
- ORB5/EUTERPE development:
 - adaptive control-variate in ORB5 (PSZS based, supported by neural networks)
 - Field-aligned finite elements in EUTERPE (collaboration with TRIMEG)
 - Further integration of the codes (already parts of a joint project; single git)
- GENE and GYSELA development and applications
 - AE + EP comparison (incl. performance) of GENE and GYSELA with PIC codes (ITPA, NL saturation, etc)
 - Realistic CoM distribution functions in GYSELA (following XTOR CoM distribution functions)
 - AE + EP + turbulence in ITER/DEMO



BACKUP SLIDES

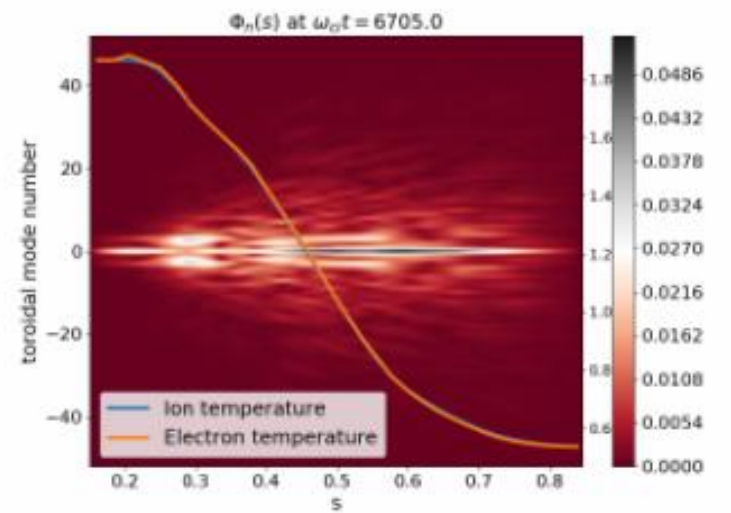
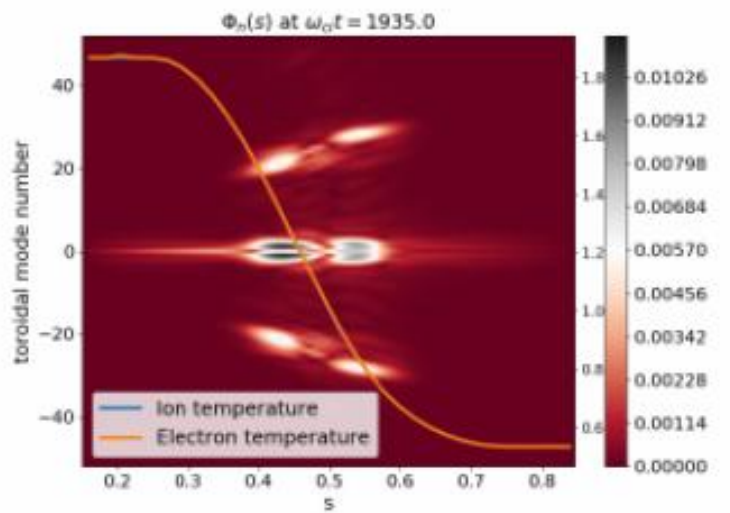
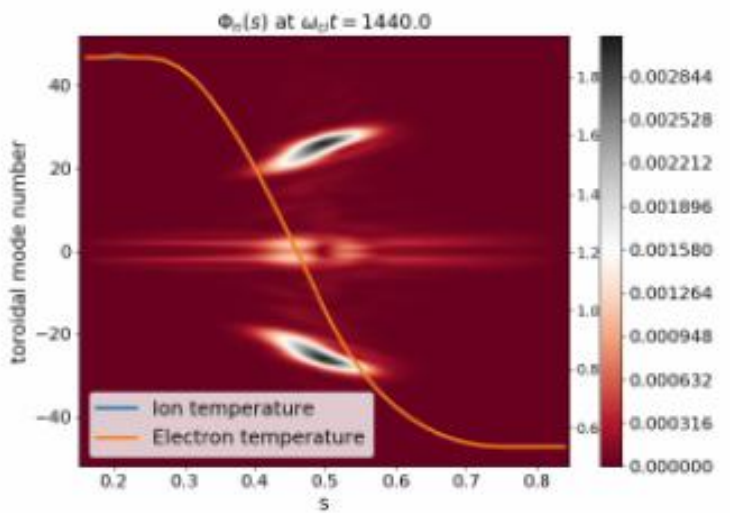
Global simulations of EM turbulence: EUTERPE



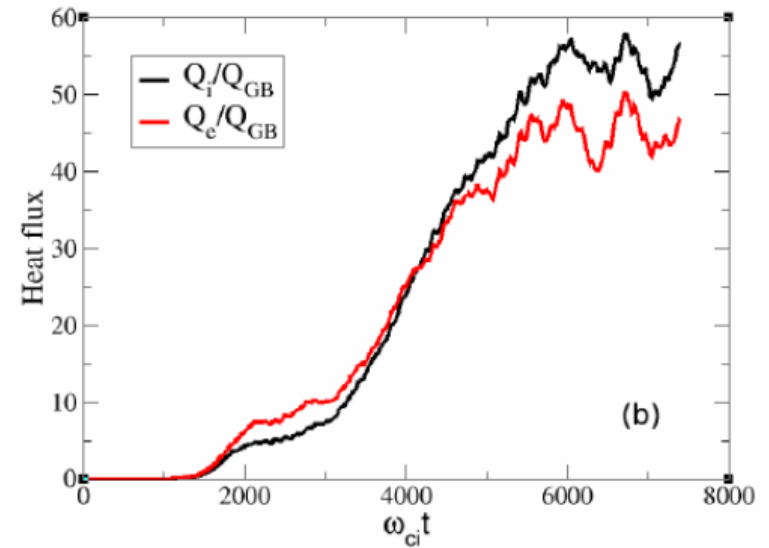
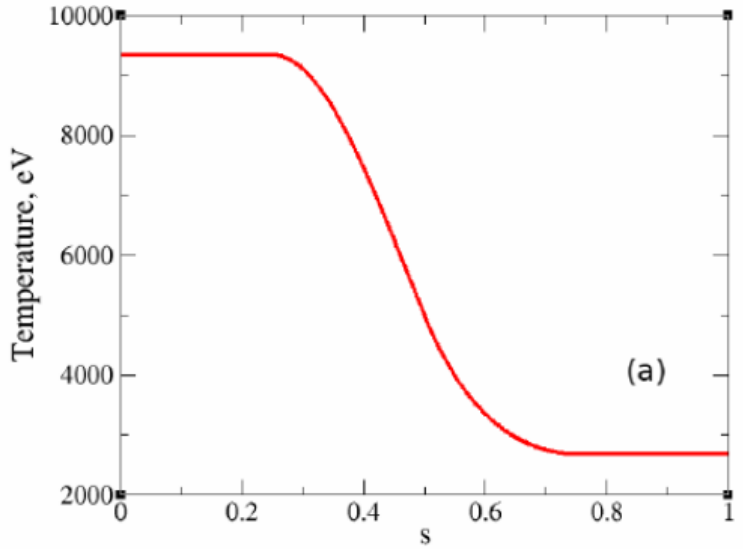
Realistic W7-X „UFM“ configuration

Global geometry is simulated
(no flux-tube issues at small $k_y \rho_i$)

For more EUTERPE results, see
poster of Y. Narbutt



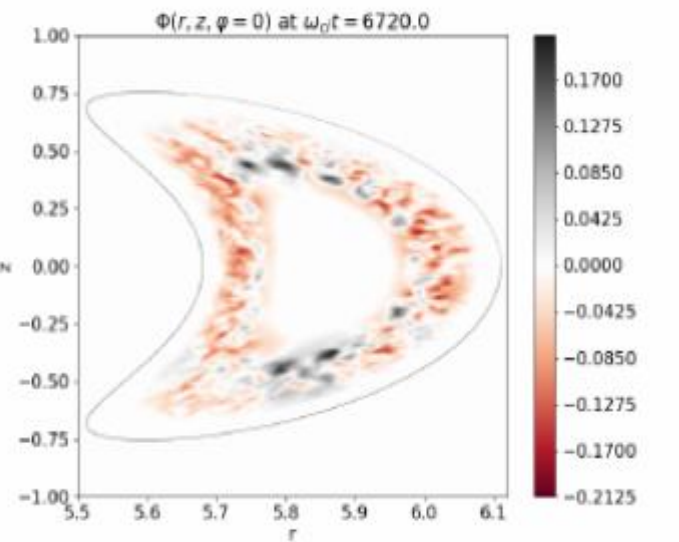
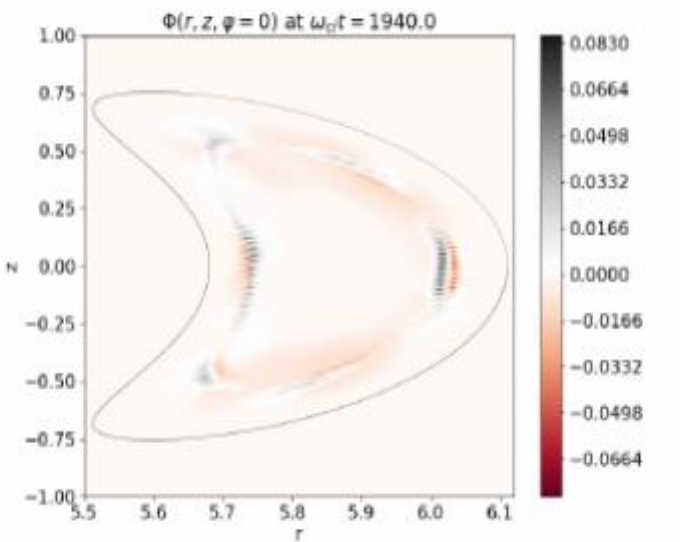
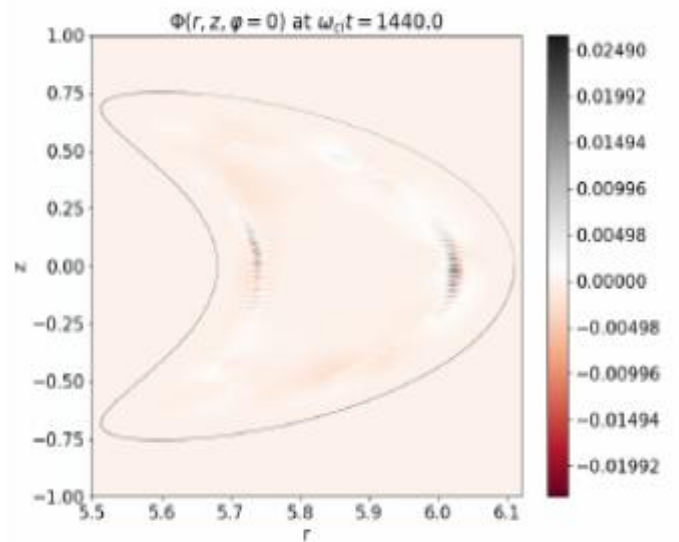
Global simulations of EM turbulence: EUTERPE



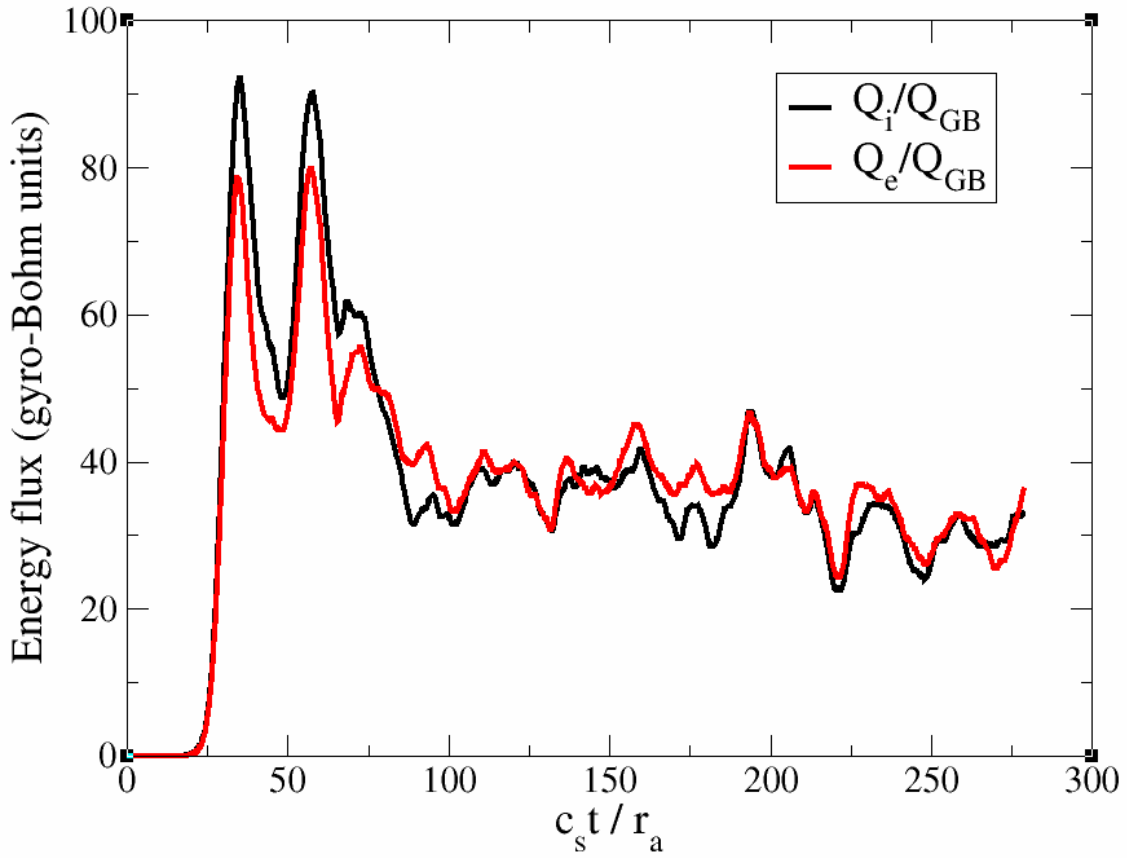
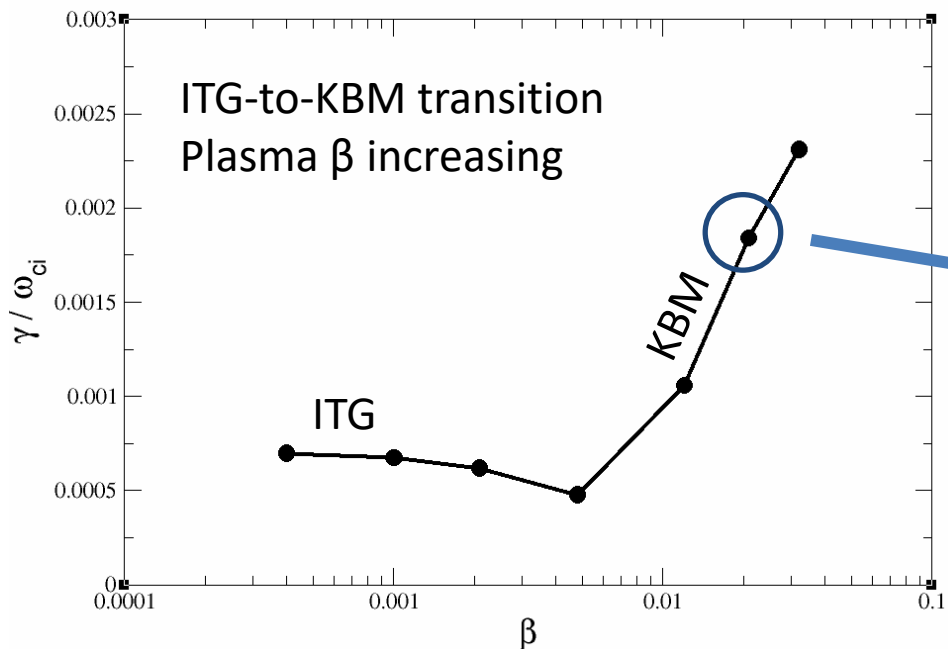
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KBM saturation and profile relaxation

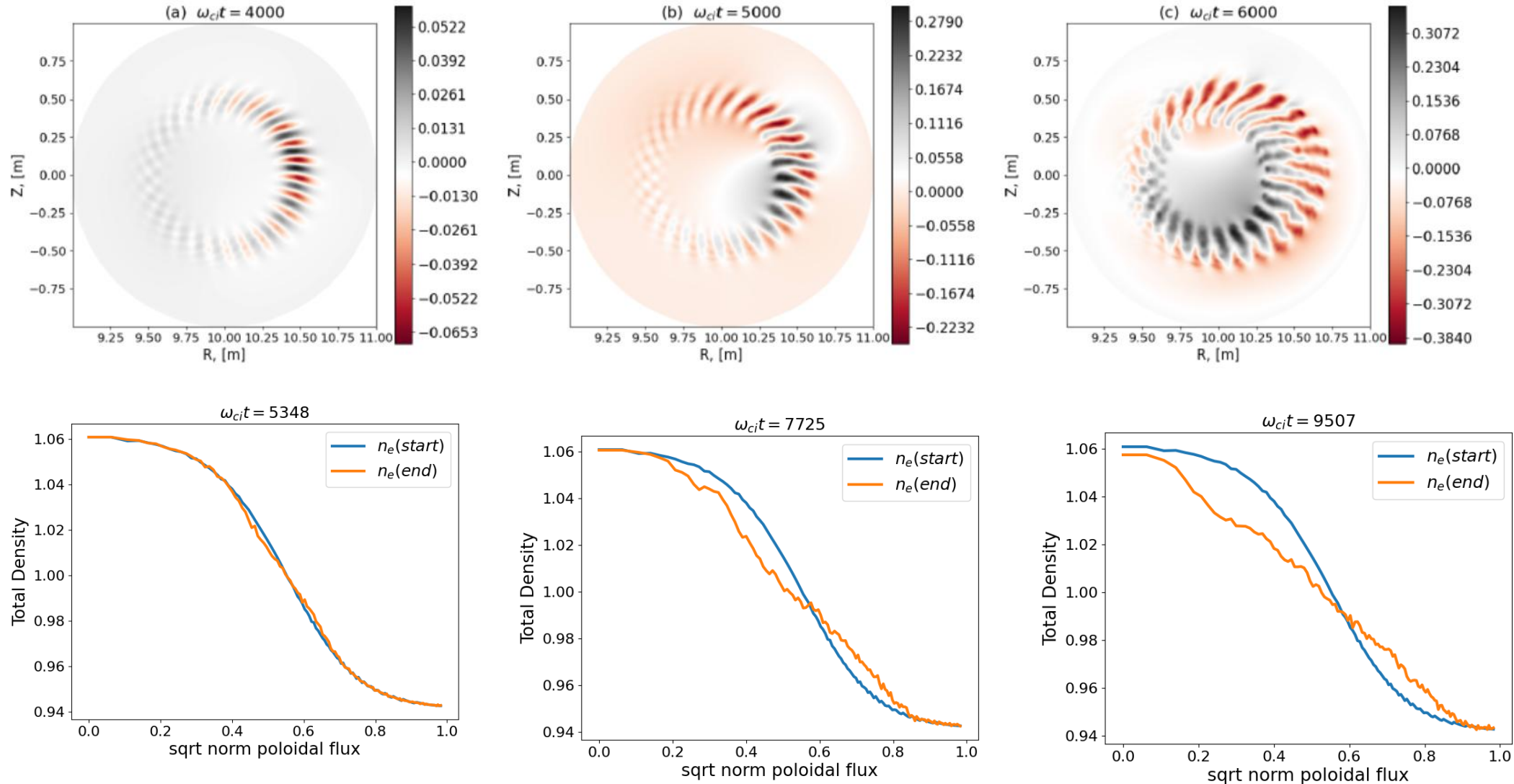


Global simulations (ORB5)

No electromagnetic run-away in global setup!

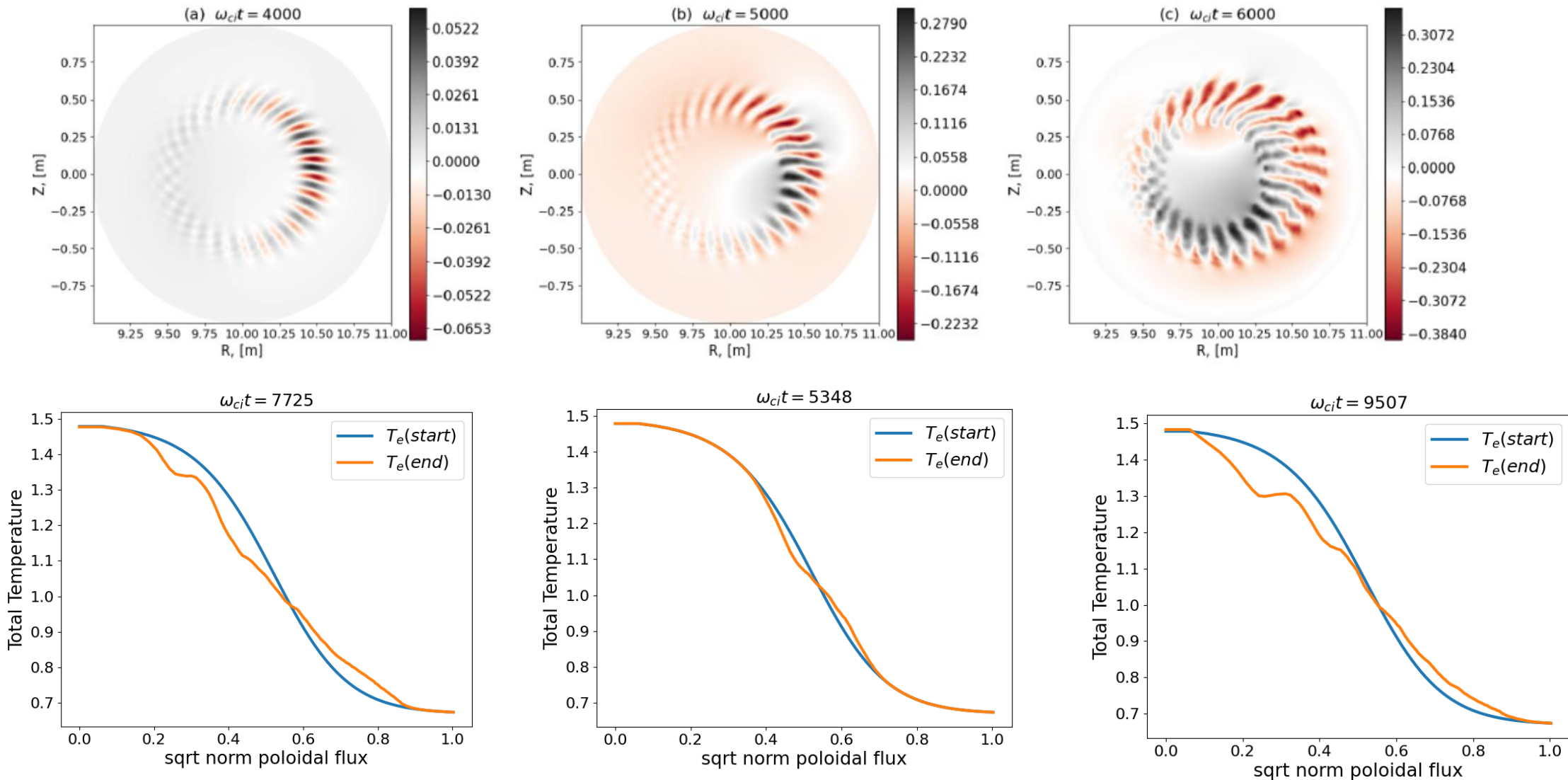


KBM saturation and profile relaxation



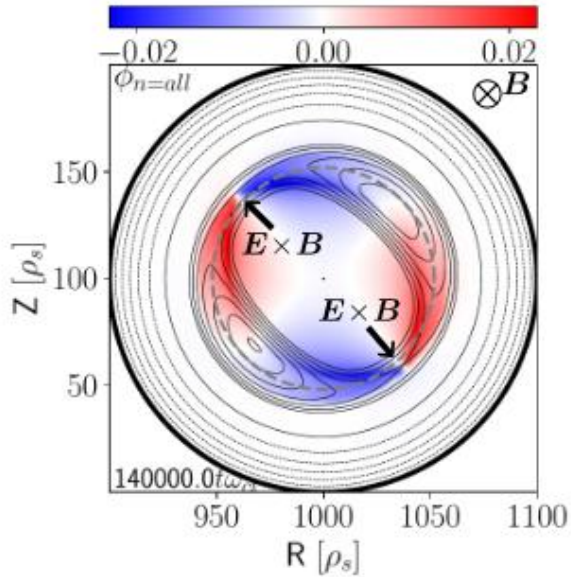
Global simulations with adaptive control variate are needed to capture the profile relaxation

KBM saturation and profile relaxation

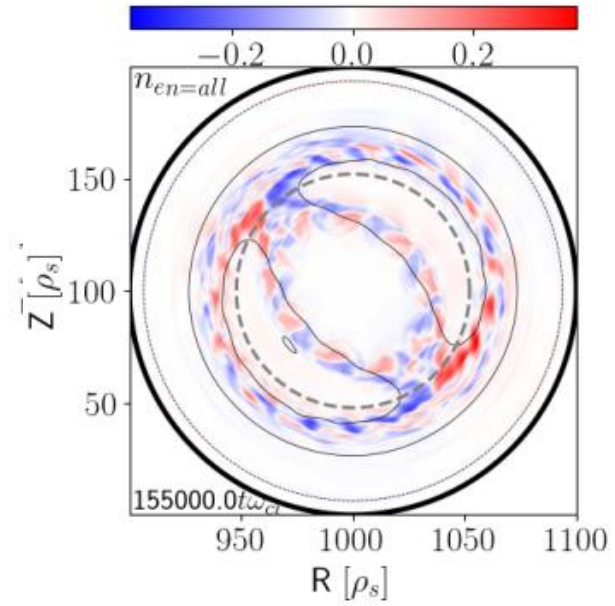


Global simulations with adaptive control variate are needed to capture the profile relaxation

Tearing instability, turbulence, effect on magnetic configuration



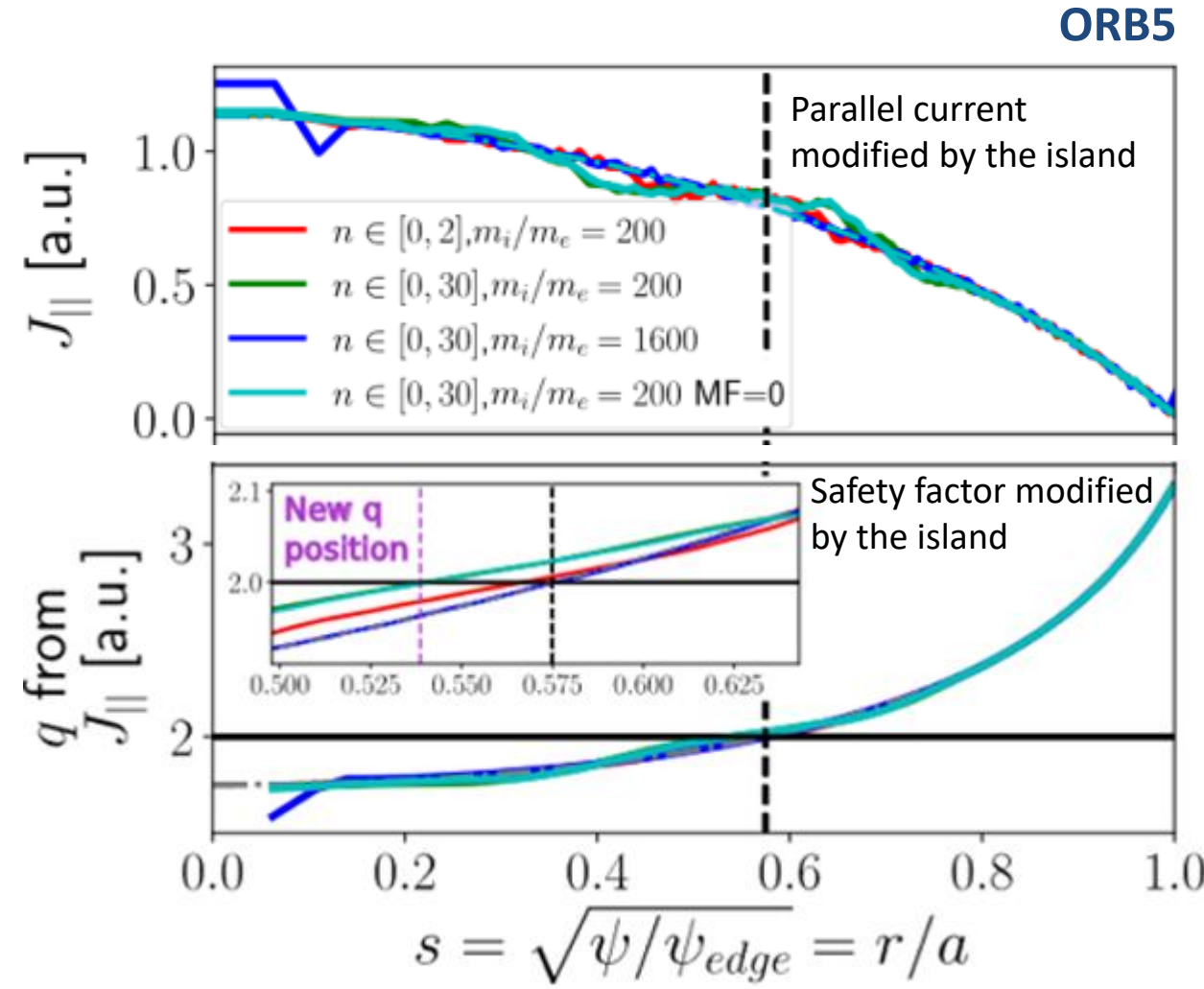
(a) Electrostatic potential ϕ .



(b) Electron density n_e .

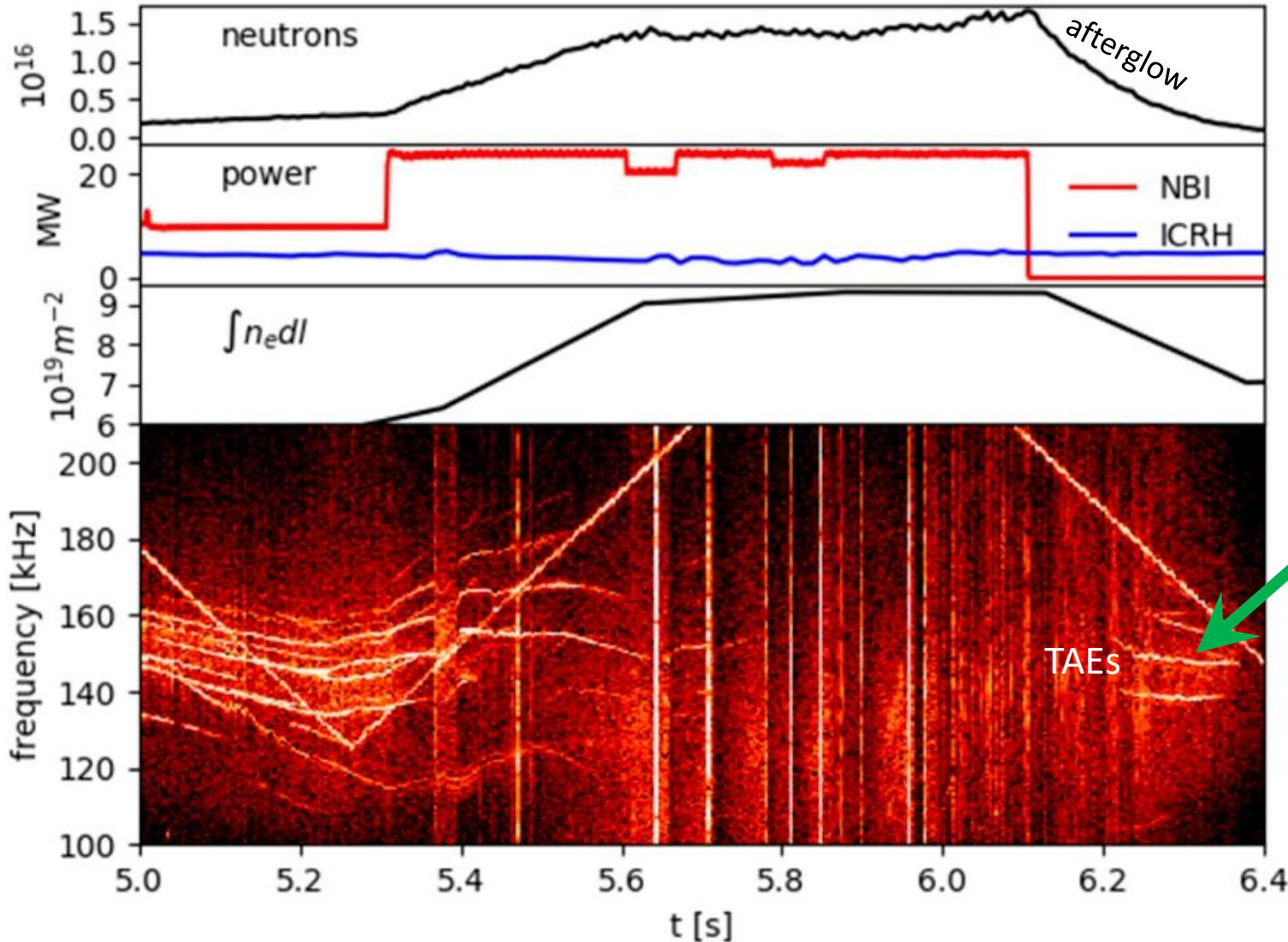
- Tearing instability generates islands**
- Islands generate flows**
- Flows become unstable (island saturation or healing).**

Tearing instability and turbulence coexist and interact
 Magnetic configuration is affected (evolution of q profile)



- Islands perturb parallel current (zonal field).**
- Safety factor and local shear nonlinearly change.**
- Island growth is nonlinearly bound.**

JET #92416: Toroidal Alfvén Eigenmodes seen in Deuterium-Deuterium experiment

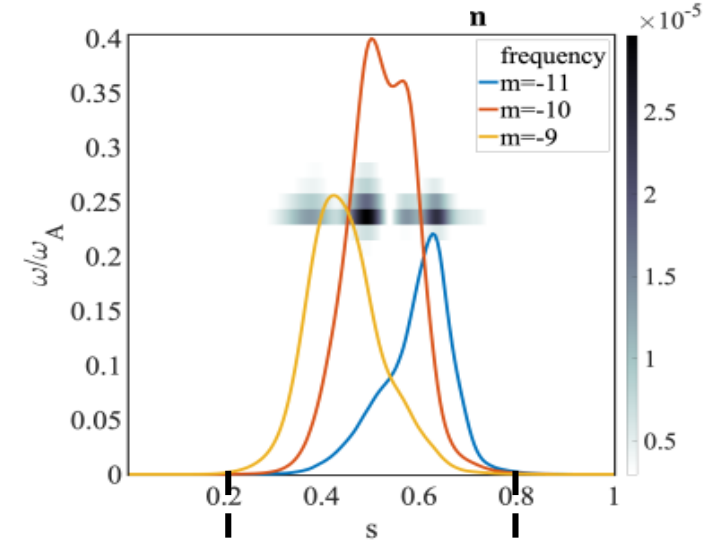
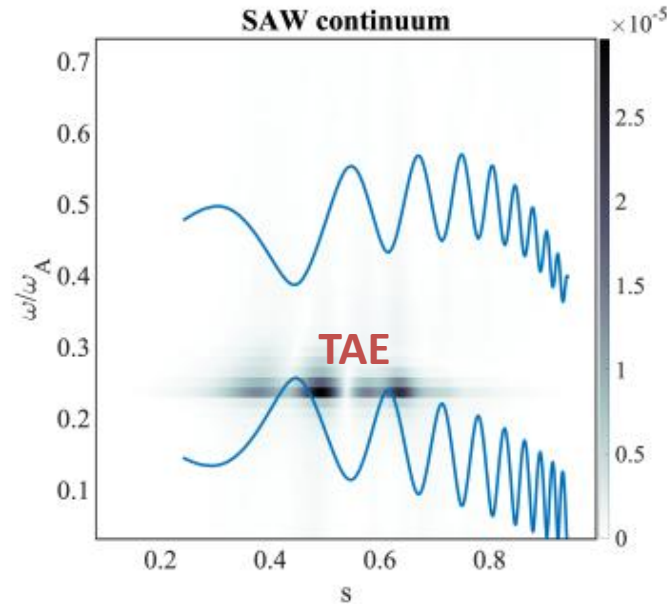
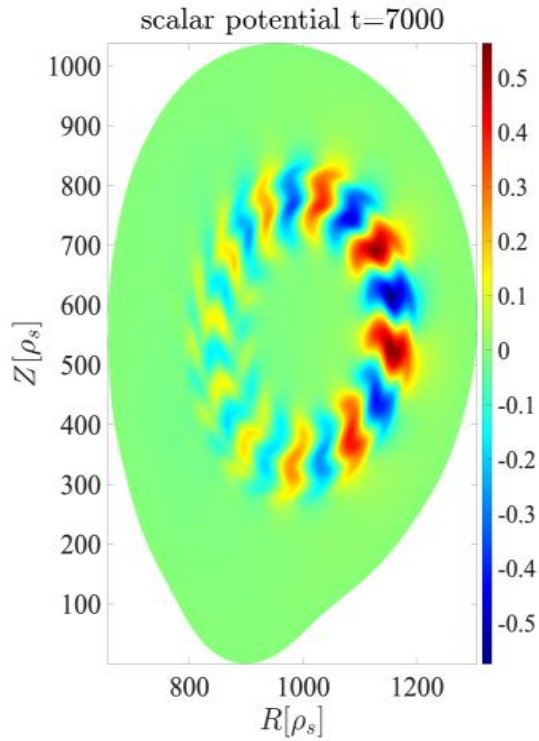


JET DD campaign, "afterglow experiment" shot #92416.

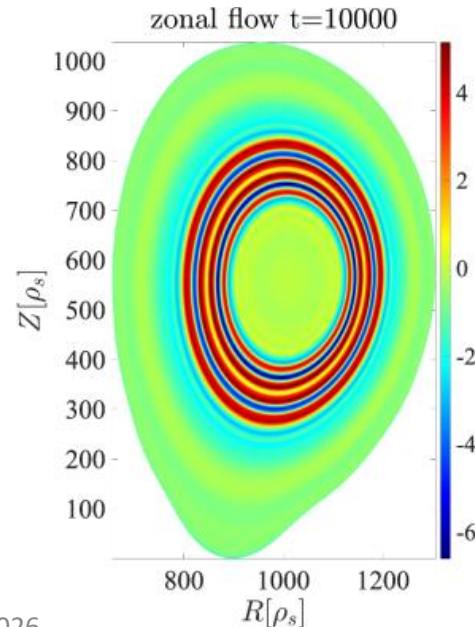
Experiment investigating plasma behavior after the main NBI heating pulse $t > 6.1$ s.

In this shot, TAEs were measured at the time 6.2 s of the discharge.

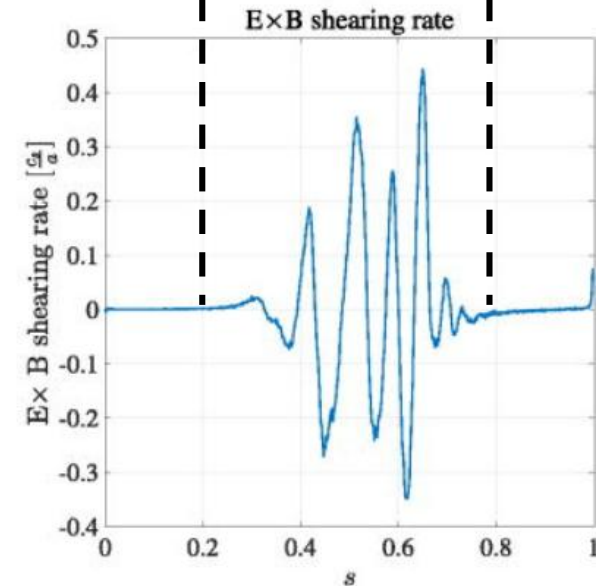
M. Fitzgerald, S.E. Sharapov, P. Siren, et al, Nuclear Fusion (2022)

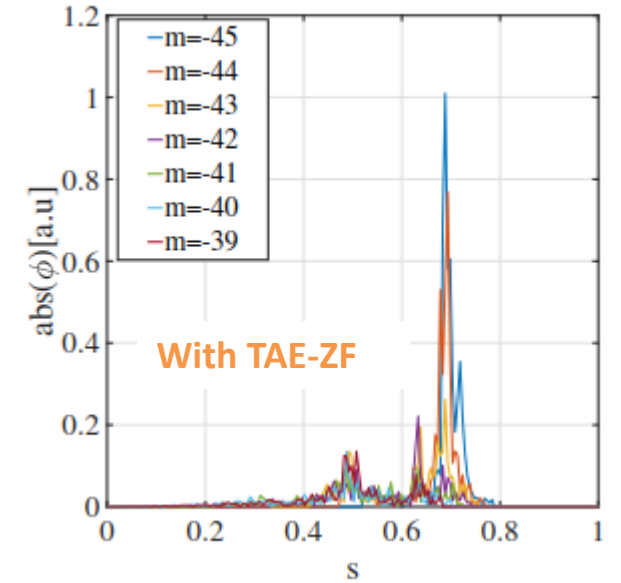
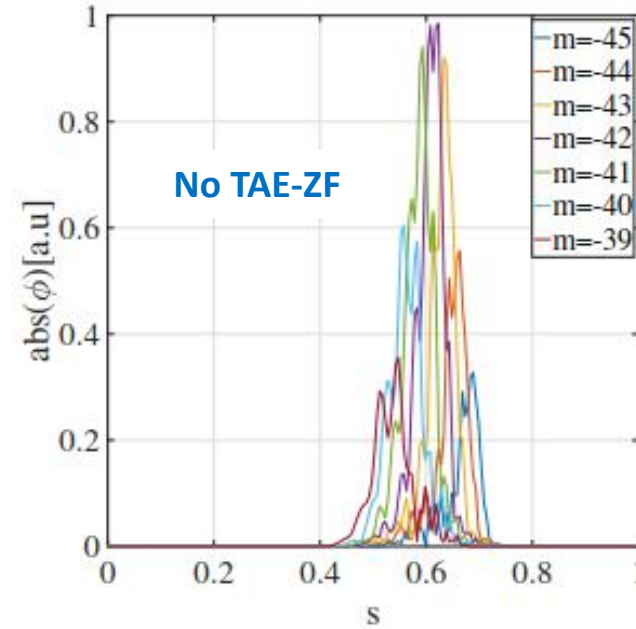
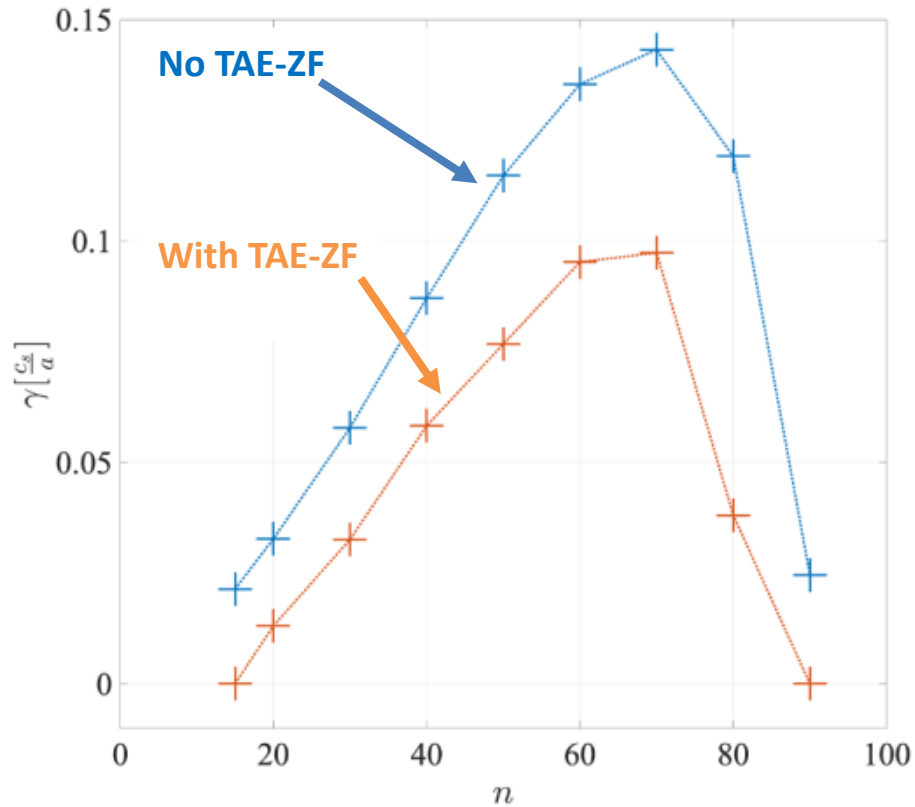


Fast-ion destabilized TAE generates beat-driven zonal flow (double growth rate of TAE).



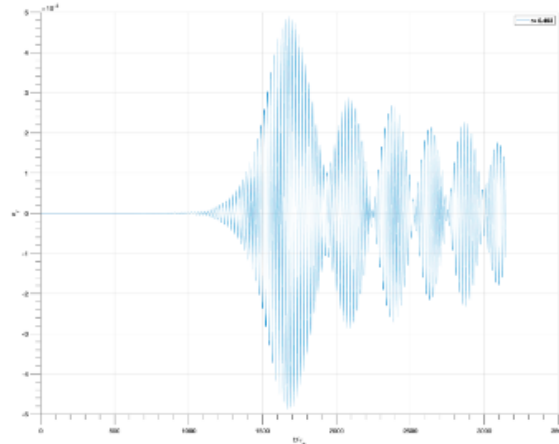
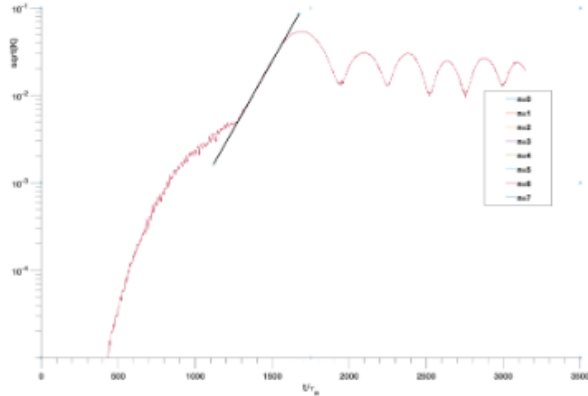
Radial location of zonal flow is determined by the TAE structure.





TAE-driven zonal flows strongly affect mode structure of ITGs

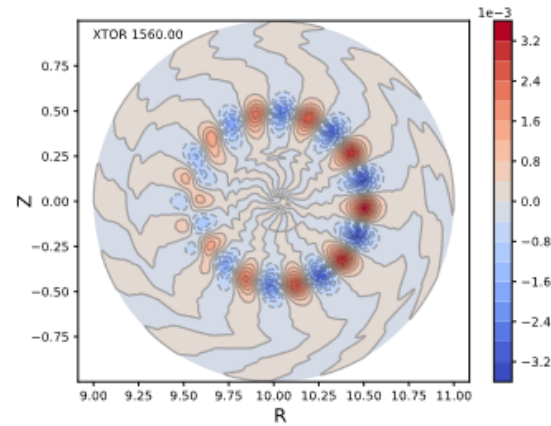
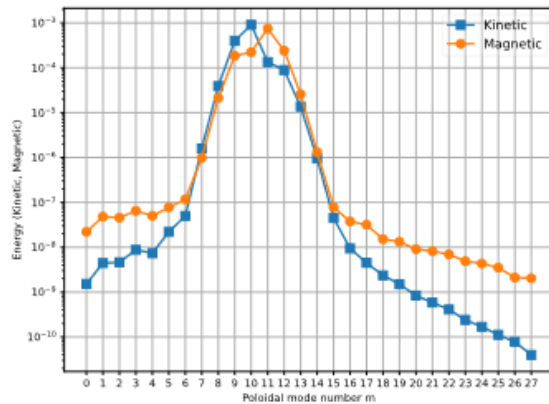
Beat-driven zonal flow mitigates ITG instability (antenna-emulated ZF)



n=6 TAE evolution:

$\Gamma = 2.18 \times 10^4 \text{ s}^{-1}$
 $\Omega = 0.399 \times 10^6 \text{ rad/s}$

Compares well with
[\[Mishchenko 2009, Könies 2018\]:](#)



$\Gamma = 2.3 \times 10^4 \pm 10\% \text{ s}^{-1}$
 $\Omega = 0.42 \times 10^6 \text{ rad/s}$

Ω ideal MHD eigenvalue
 code (CAS3D):
 $\Omega = 0.401 \times 10^6 \text{ rad/s}$



Call for TSVV-G: Physics of Burning Plasmas

Develop a self-consistent description and corresponding simulation tools for the mutual interaction of energetic particles, MHD modes, turbulence, and kinetic plasma profiles in both tokamak and stellarator geometries.

Develop a theoretical understanding and validated interpretative/predictive capabilities (including reduced models) of burning plasma physics in both tokamak and stellarator geometries. The resulting reduced models are to be used in TSVV-H and TSVV-I.

Develop strategies to optimize the deposition of fusion alpha energy into the bulk plasma to enhance reactor performance.